

Velocity-Stress Finite-Difference Modeling of Poroelastic Wave Propagation

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and

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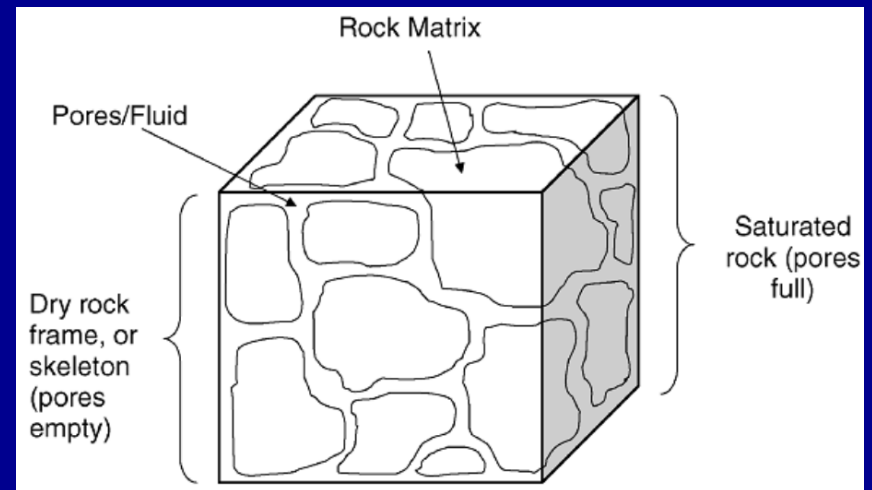


Outline

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- Biot's Theory
- Staggered-Grid Finite Difference
- Numerical Examples
- Conclusion
- Acknowledgement

Introduction

- Poroelastic Medium



(Russell et al., 2003)

- Biot (1962): anelastic effects from the relative movement of the fluid.
- Biot's theory: Important in oil and gas exploration, CO₂ storage monitoring and hydrogeology.
- The Theory predicts two compressional waves and one shear wave.

Biot's Theory(1962)

Assumptions :

- Elastic rock frame
- Connected pores
- Seismic wavelength \gg average pore size
- Small deformations
- Statistically isotropic medium

- Stress-Strain Relation For Porous Media (Biot, 1962)

Solid Stress $\tau_{ij} = 2\mu e_{ij} + (\lambda_c e_{kk} + \alpha M \varepsilon_{kk}) \delta_{ij}$

Fluid Pressure $P = -\alpha M e_{kk} - M \varepsilon_{kk}$

$$e_{ij} = \nabla \cdot u = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

$$\varepsilon_{ij} = \nabla \cdot (u - U)$$

u : Solid Particle Displacement

U : Fluid Particle Displacement

$$\alpha = 1 - \frac{K_{Dry}}{K_{Solid}}$$

$$M = \left[\frac{\phi}{K_{Fluid}} + \frac{(\alpha - \phi)}{K_{Solid}} \right]$$

Coupling Modulus

λ & μ : Lamé Parameters of the Saturated Rock.

- Equations of motion for a statistically isotropic porous media saturated with viscous fluid:

$$(m\rho - \rho_f^2) \frac{\partial^2 u_i}{\partial t^2} = m \frac{\partial \tau_{ij}}{\partial x_j} + \rho_f b \frac{\partial w_i}{\partial t} + \rho_f \frac{\partial P}{\partial x_i}$$

$$(m\rho - \rho_f^2) \frac{\partial^2 w_i}{\partial t^2} = -\rho_f \frac{\partial \tau_{ij}}{\partial x_j} - \rho b \frac{\partial w_i}{\partial t} - \rho \frac{\partial P}{\partial x_i}$$

Effective
Fluid Density

$$m = T \frac{\rho_f}{\phi}$$

ρ_f : Fluid Density

ρ : Density of Saturated
Rock

Fluid Displacement
Relative to the Solid

$$w = u - U$$

Mobility
 $b = \eta/\kappa$

η : Viscosity

κ : Permeability

Substituting $V = \frac{\partial u}{\partial t}$ and $W = \frac{\partial w}{\partial t}$ in the equations of motion and taking derivatives with respect to time from both sides of the stress-strain relationship we have:

$$\begin{aligned}(m\rho - \rho_f^2) \frac{\partial V_i}{\partial t} &= m \frac{\partial \tau_{ij}}{\partial x_j} + \rho_f b W + \rho_f \frac{\partial P}{\partial x_i} \\(m\rho - \rho_f^2) \frac{\partial W_i}{\partial t} &= -\rho_f \frac{\partial \tau_{ij}}{\partial x_j} - \rho b W - \rho \frac{\partial P}{\partial x_i}\end{aligned}$$

and

$$\frac{\partial \tau_{ij}}{\partial t} = 2\mu \frac{\partial e_{ij}}{\partial t} + \left(\lambda_c \frac{\partial e_{kk}}{\partial t} + \alpha M \frac{\partial \varepsilon_{kk}}{\partial t} \right) \delta_{ij}$$

$$\frac{\partial P}{\partial t} = -\alpha M \frac{\partial e_{kk}}{\partial t} - M \frac{\partial \varepsilon_{kk}}{\partial t}$$

- 2D case:

$$\frac{\partial \tau_{xx}}{\partial t} = (\lambda_c + 2\mu) \frac{\partial V_x}{\partial x} + \lambda_c \left(\frac{\partial V_z}{\partial z} \right) + \alpha M \left(\frac{\partial W_x}{\partial x} + \frac{\partial W_z}{\partial z} \right) \quad (1)$$

$$\frac{\partial \tau_{zz}}{\partial t} = (\lambda_c + 2\mu) \frac{\partial V_z}{\partial z} + \lambda_c \left(\frac{\partial V_x}{\partial x} \right) + \alpha M \left(\frac{\partial W_x}{\partial x} + \frac{\partial W_z}{\partial z} \right) \quad (2)$$

$$\frac{\partial \tau_{xz}}{\partial t} = \mu \left(\frac{\partial V_z}{\partial x} + \frac{\partial V_x}{\partial z} \right) \quad (3)$$

$$\frac{\partial P}{\partial t} = -\alpha M \left(\frac{\partial V_x}{\partial x} + \frac{\partial V_z}{\partial z} \right) - M \left(\frac{\partial W_x}{\partial x} + \frac{\partial W_z}{\partial z} \right) \quad (4)$$

$$\frac{\partial V_x}{\partial t} = A \left(\frac{\partial \tau_{xx}}{\partial x} + \frac{\partial \tau_{xz}}{\partial z} \right) + BW_x + C \frac{\partial P}{\partial x} \quad (5)$$

$$\frac{\partial V_z}{\partial t} = A \left(\frac{\partial \tau_{zx}}{\partial x} + \frac{\partial \tau_{zz}}{\partial z} \right) + BW_z + C \frac{\partial P}{\partial z} \quad (6)$$

$$\frac{\partial W_x}{\partial t} = D \left(\frac{\partial \tau_{xx}}{\partial x} + \frac{\partial \tau_{xz}}{\partial z} \right) + EW_x + F \frac{\partial P}{\partial x} \quad (7)$$

$$\frac{\partial W_z}{\partial t} = D \left(\frac{\partial \tau_{zx}}{\partial x} + \frac{\partial \tau_{zz}}{\partial z} \right) + EW_z + F \frac{\partial P}{\partial z} \quad (8)$$

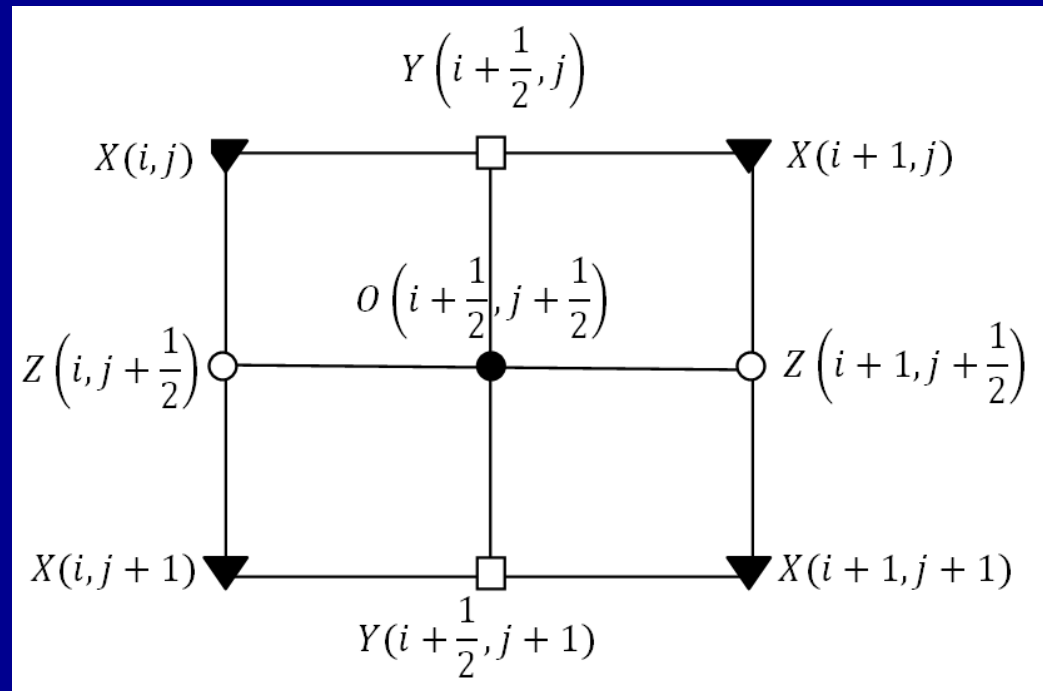
Staggered-Grid Finite Difference(Levander, 1988)

X : τ_{xx}, τ_{zz} and P

Y : V_x and W_x

Z : V_z and W_z

O : τ_{xz}

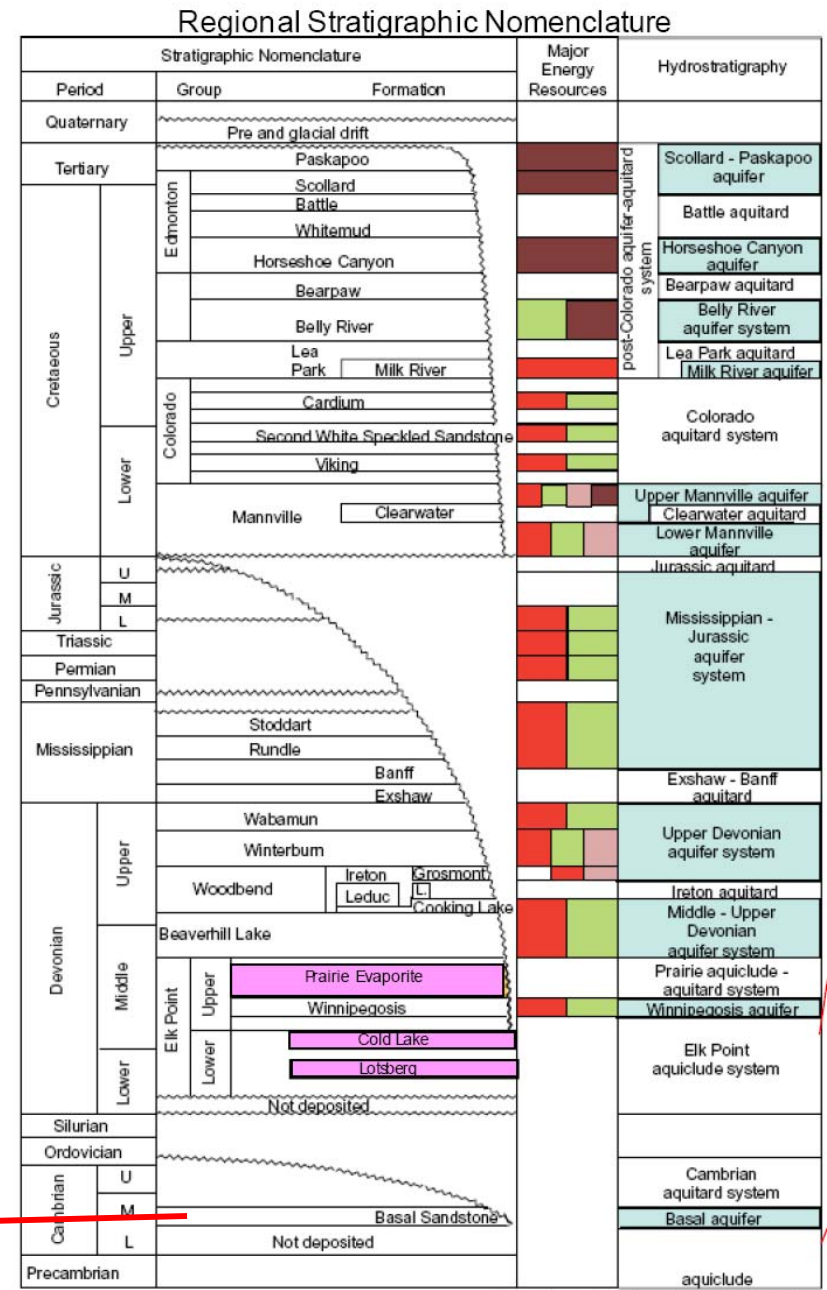


Numerical Examples

Single layer model based on QUEST Project

- CO₂ storage in Basal Cambrian Sands or BCS, which is a saline aquifer within Western Canadian Sedimentary Basin (WCSB)
- Data from well SCL-8-19-59-20W4

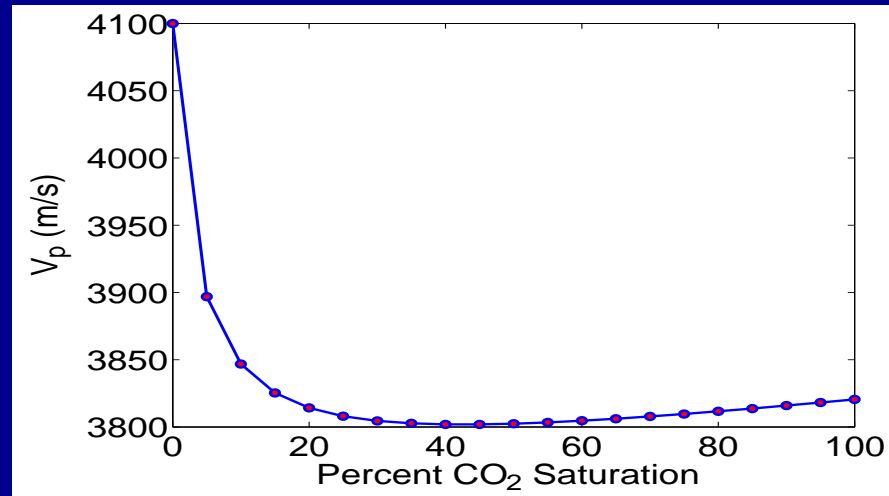
Quest



Modified after Bachu et al. 2000.

Single Layer Model

Gassmann Fluid Substitution



ρ_f	937 (kg/m^3)
ρ	2370 (kg/m^3)
V_p	3800 (m/s)
V_s	2400 (m/s)
ϕ	16%
κ	1(mD)

BCS: 40% CO_2

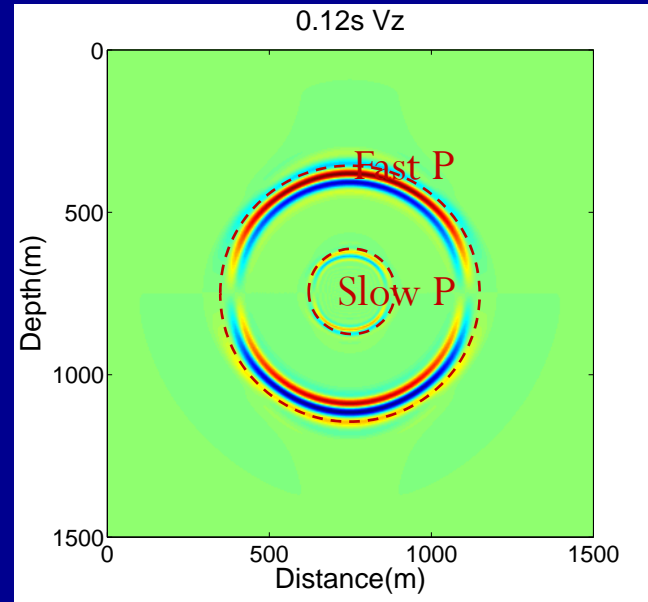
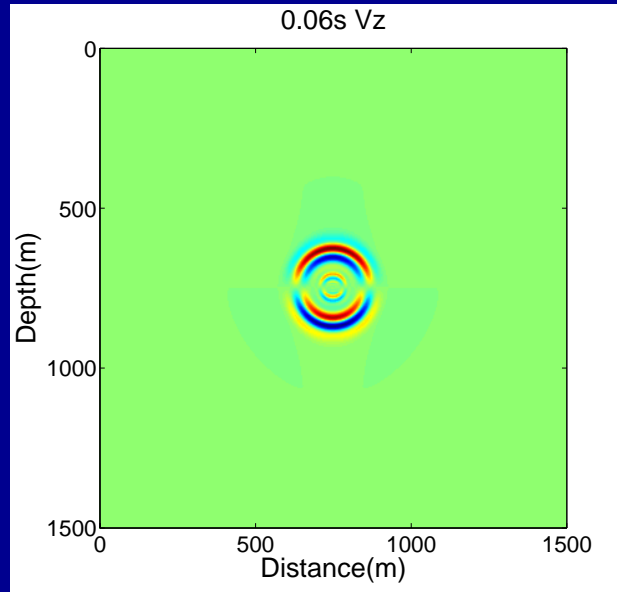
- Fourth order in space and second order in time.
- The stability condition is the same as the one in the elastic case (Zhu:1991)

$$\Delta t \leq \frac{h}{(V_p^2 - V_s^2)^{1/2}}$$

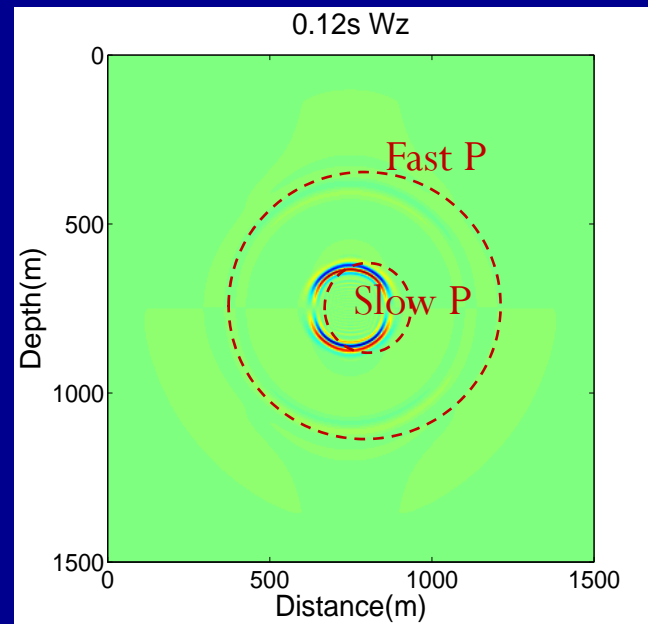
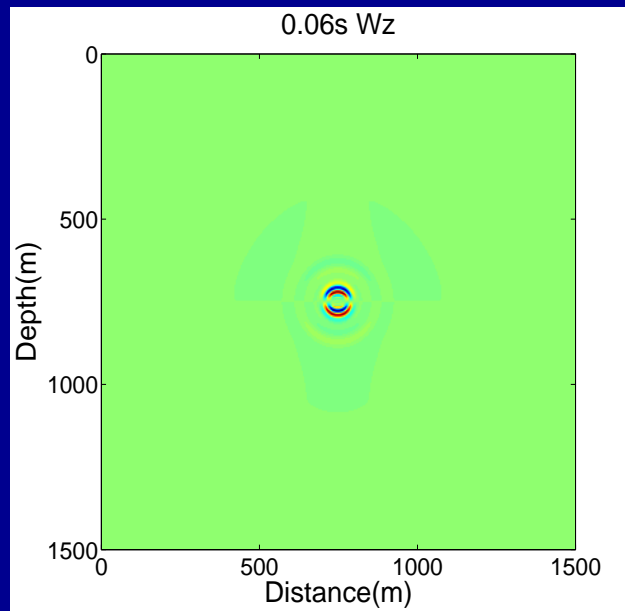
$$h = 3m \quad dt = 0.2 \text{ ms}$$

- The size of the model was 1500 m by 1500 m
- Explosive source: Ricker wavelet with dominant frequency 50 Hz
- Source location : $(x, z) = (750, 750)m$

- Vertical Particle Velocity Snapshots:

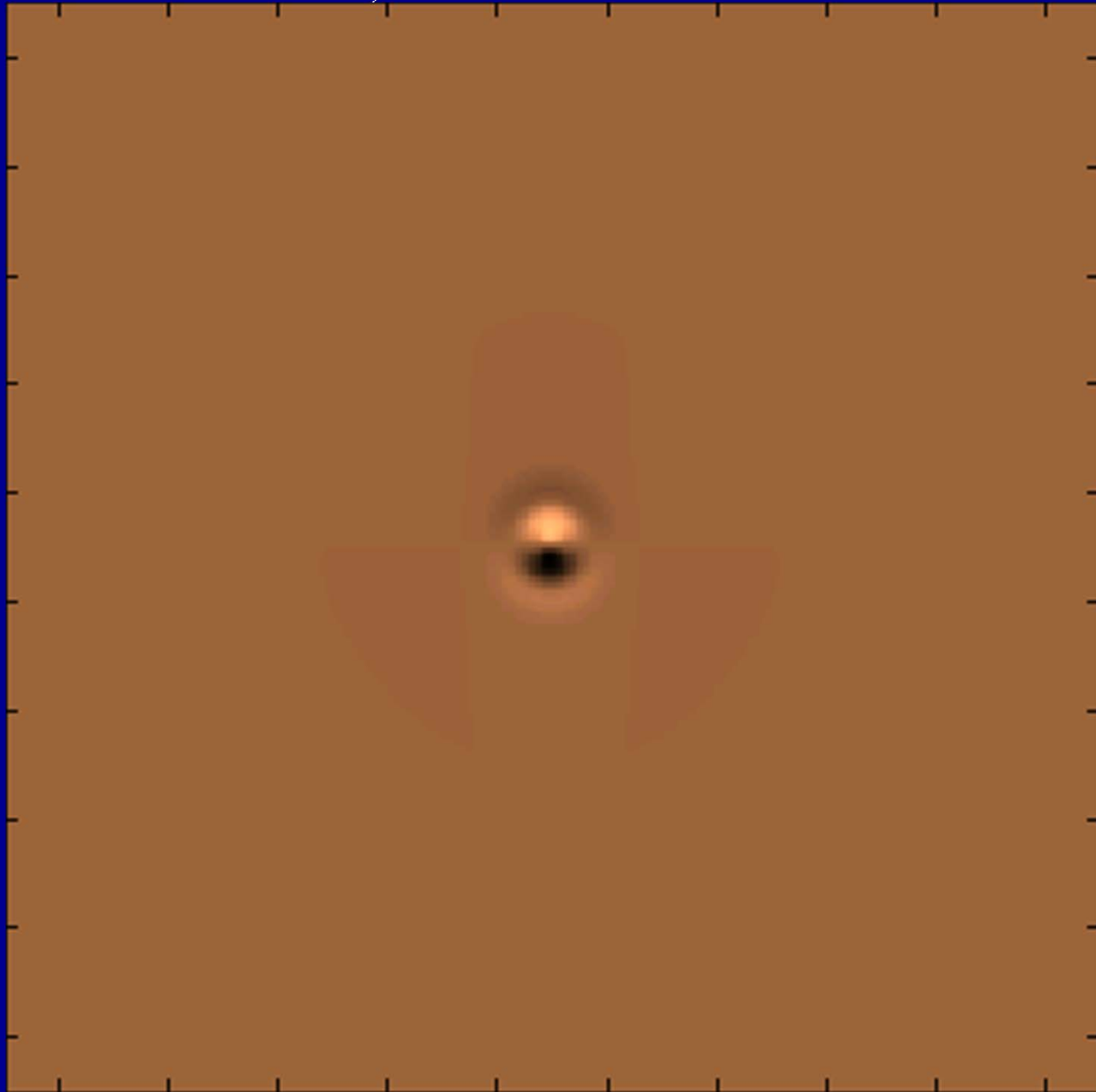


Solid

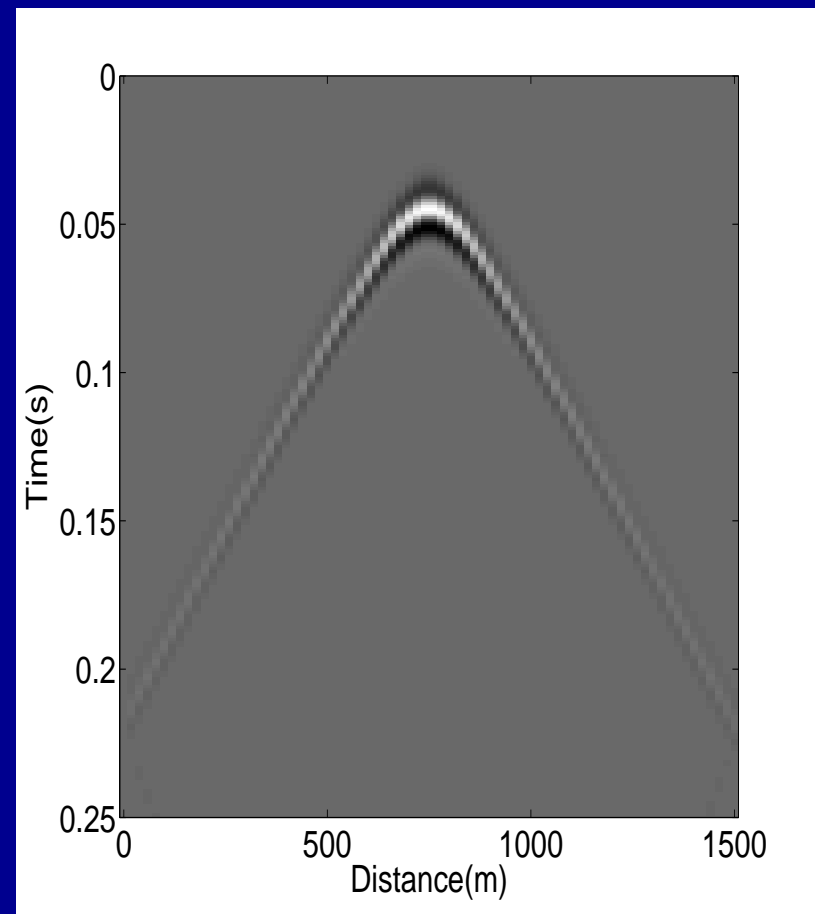
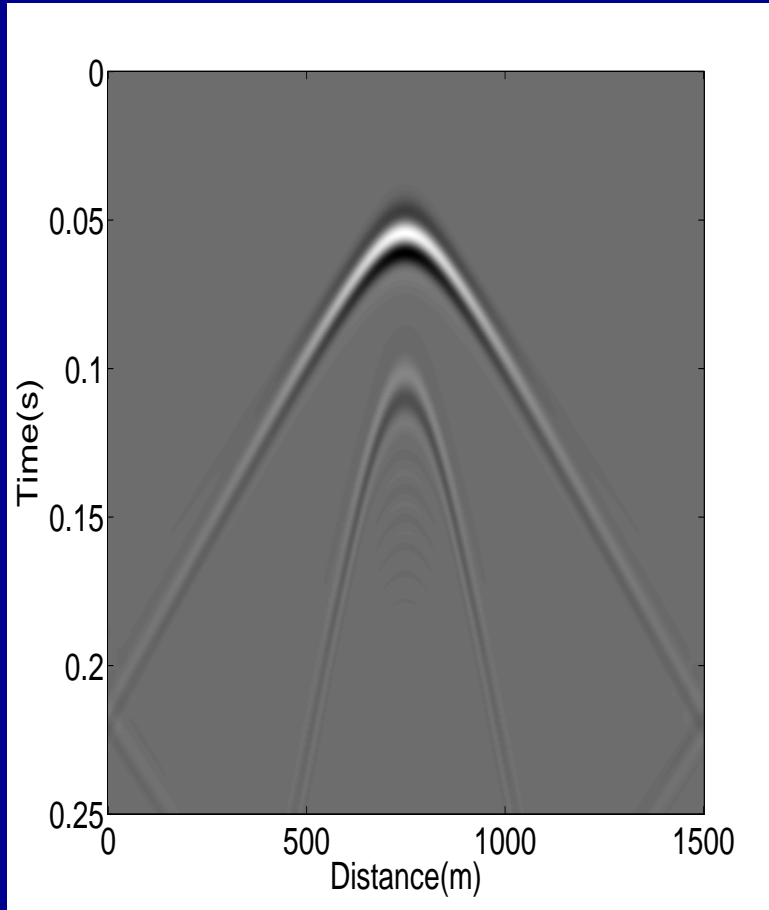


Fluid

- Vertical Particle Velocity of the Solid

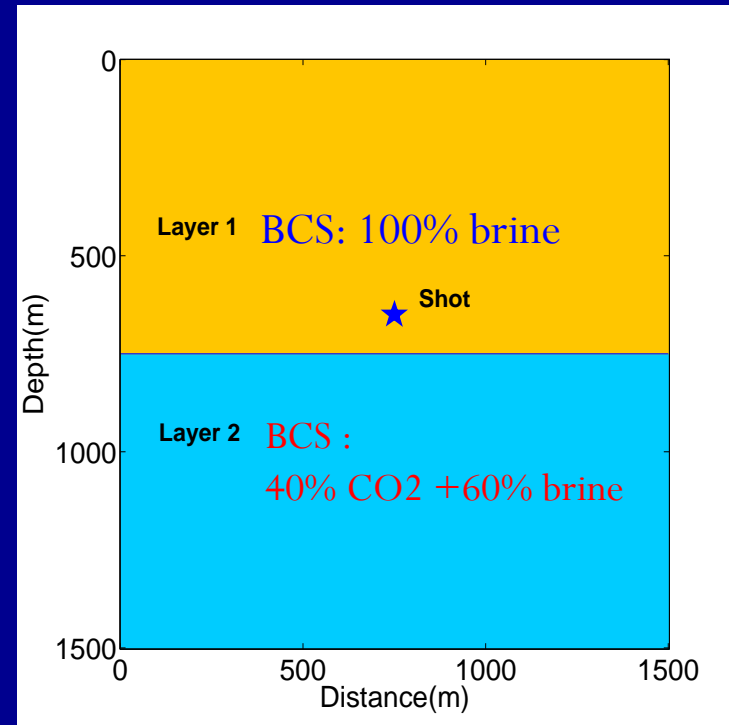


- Comparison with elastic algorithm

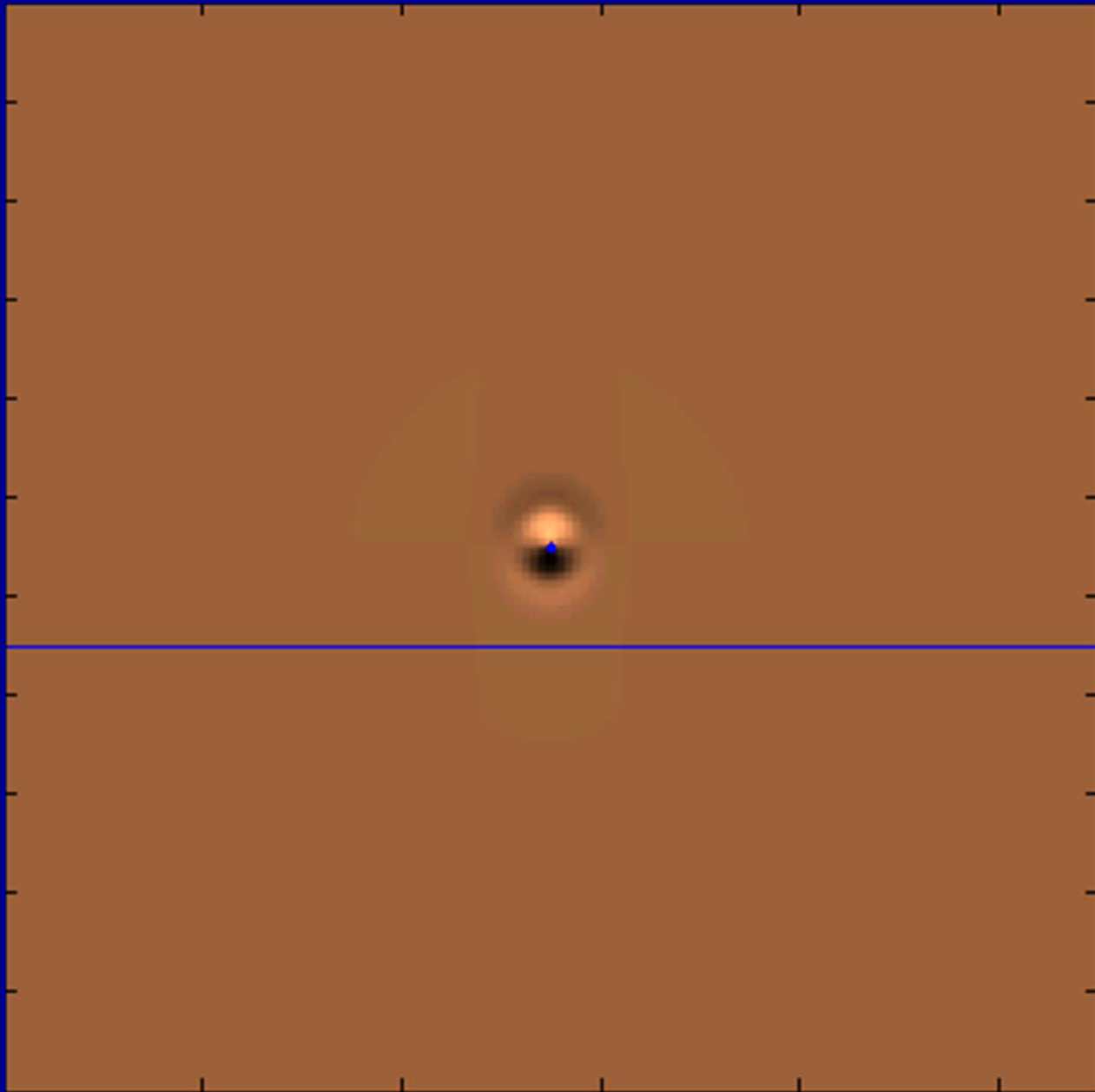


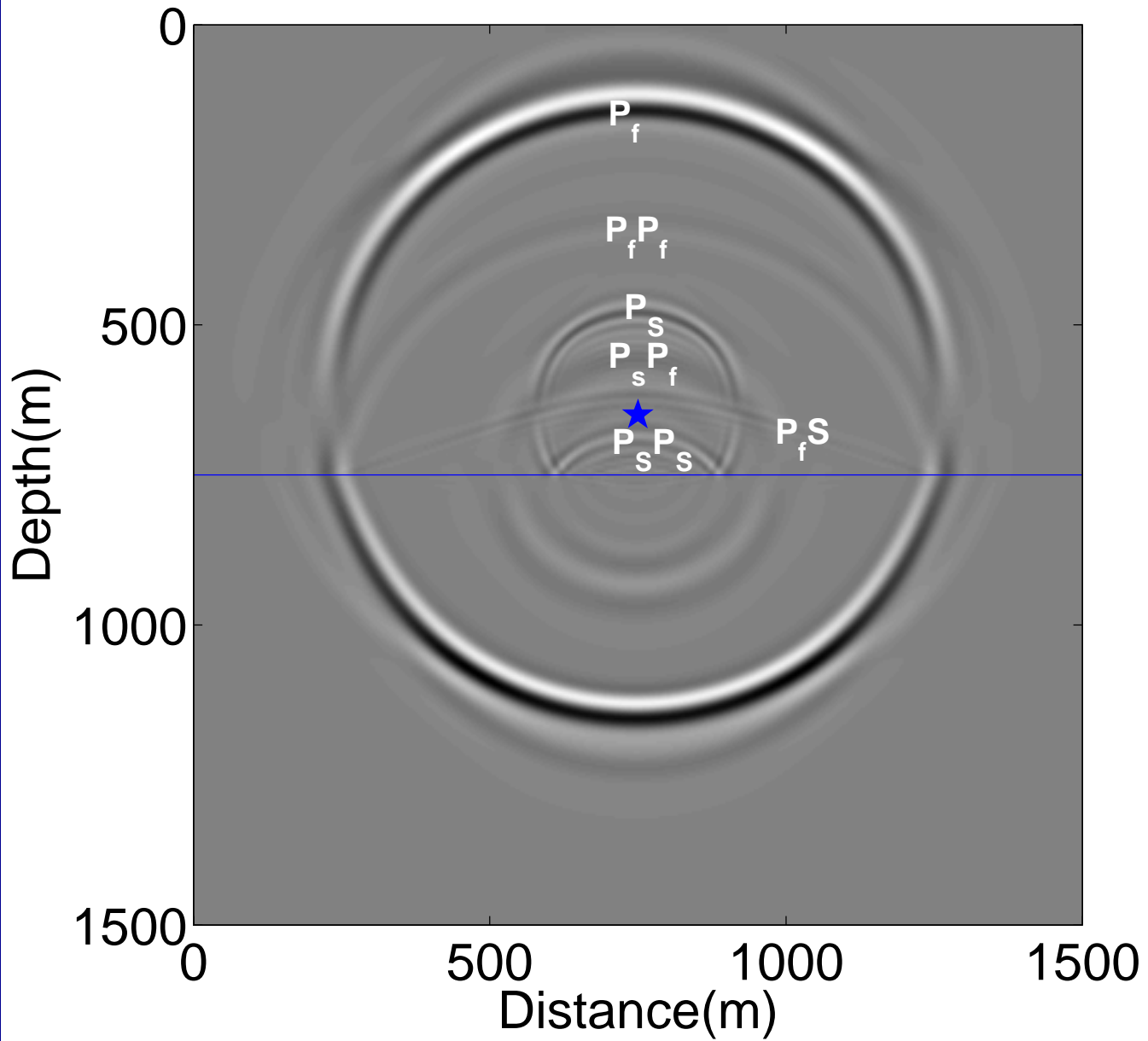
Two-Layered Model

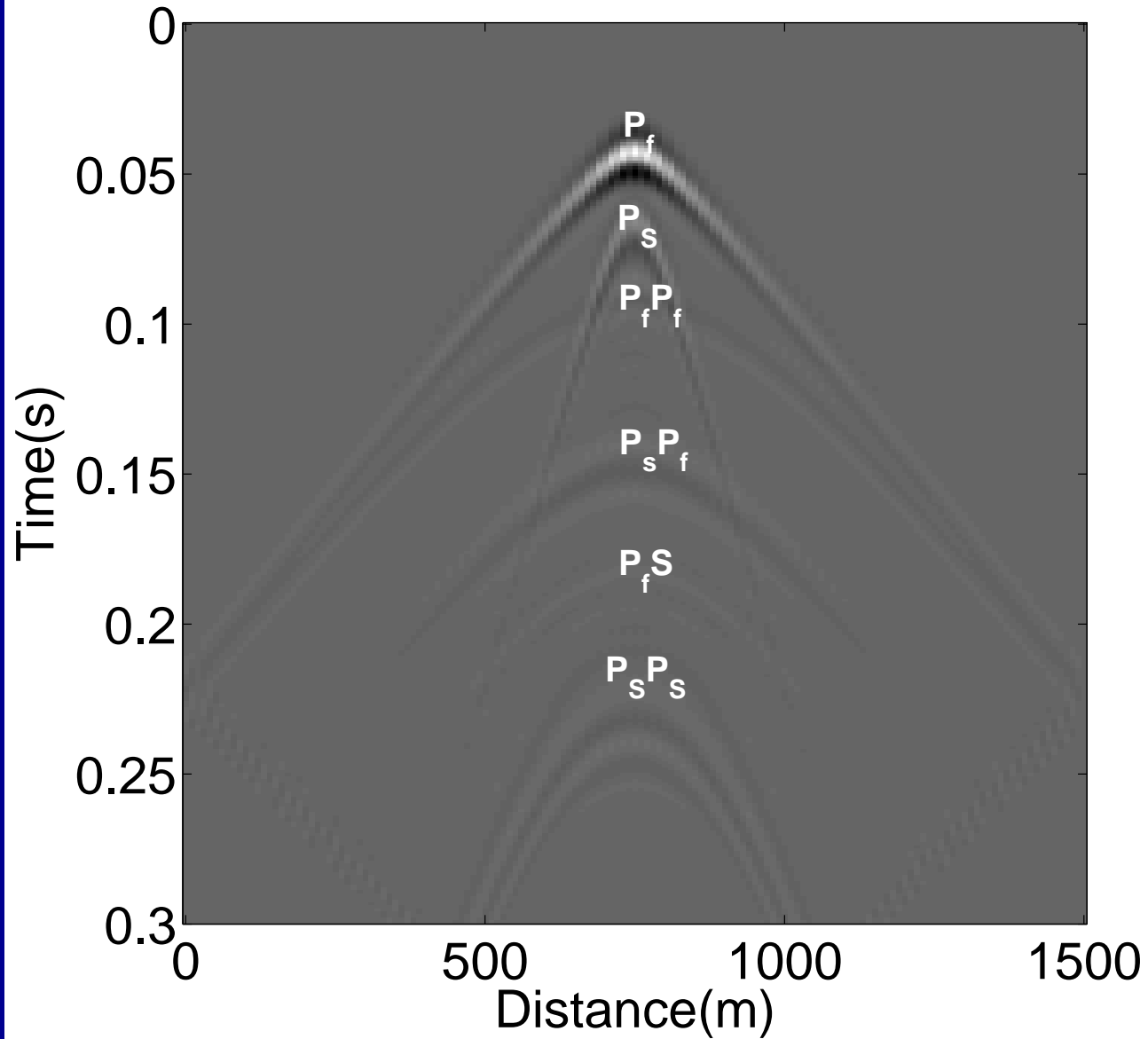
	Top Layer	Bottom Layer
ρ_f	1070 (kg/m^3)	937 (kg/m^3)
ρ	2400 (kg/m^3)	2370 (kg/m^3)
V_p	4100(m/s)	3800 (m/s)
V_s	2390(m/s)	2400 (m/s)
ϕ	16%	16%
κ	1(mD)	1(mD)



- Vertical Particle Velocity of the Solid







Conclusion and Future Goals

- The Poroelastic algorithm Generates slow compressional wave as predicted by Biot's theory.
- At a poroelastic boundary the slow P-wave is converted to a fast P-wave.
- The algorithm handles layered models and should be examined for more complex models.
- The algorithm could be used for inversion to obtain porous media properties that are ignored in elastic algorithms.

Thanks to:

- Carbon Management Canada (CMC) for financial support.
- CREWES project for extensive technical support.
- Hassan Khaniani, Peter Manning and Joe Wong from CREWES
- David Aldridge from Sandia National Laboratories.
- Shell for the data

THANKS!

- Fluid Pressure

