

Circular Wavefront Assumptions for Gridded Traveltime Computations

John C. Bancroft

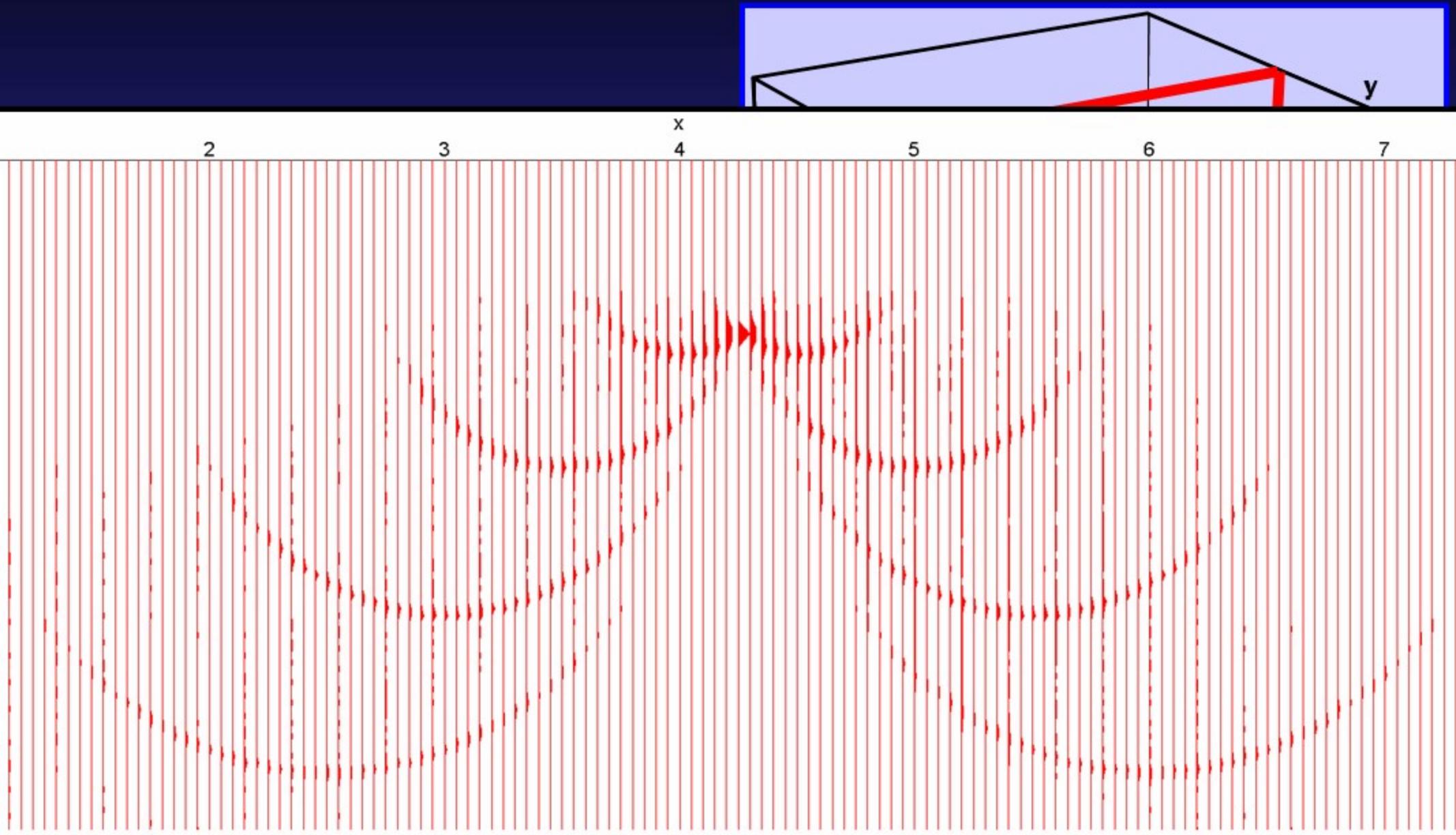
CREWES/University of Calgary

CREWES 2005

Synopsis

- GPR
- Fast-filtering and the Gabor transform
- Multigrid... statics
- Circular wavefront
- Anisotropy

GPR 3D migration of 2D gridded data

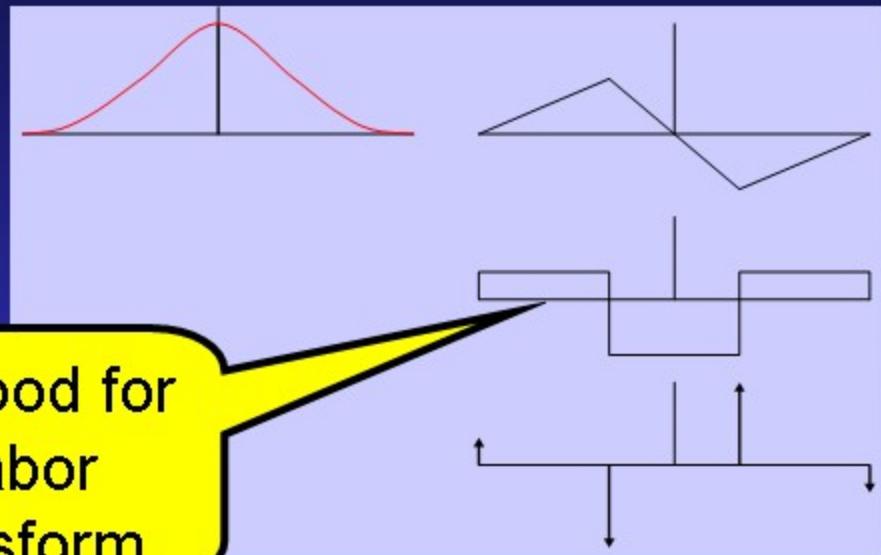


Fast-filtering...

$$c(t) = \int_{-\infty}^{\infty} w(t-\tau)g(\tau)d\tau \quad C(f) = W(f)G(f)$$

$$C(f) = [j\omega]^2 W(f) \frac{1}{[j\omega]^2} G(f)$$

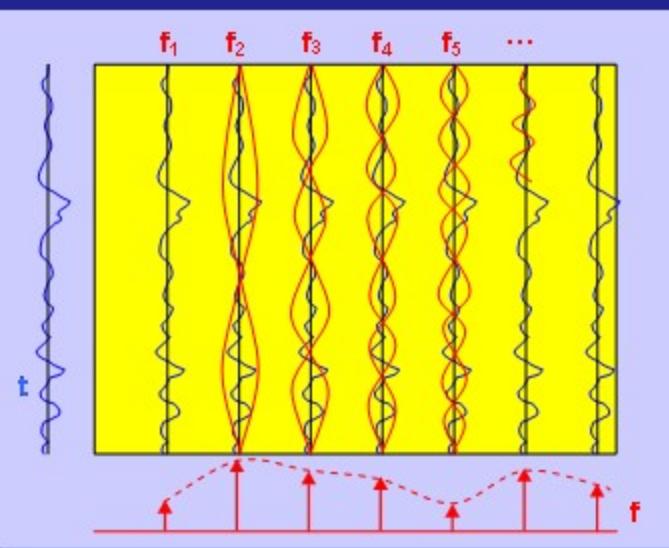
Piecewise quadratic



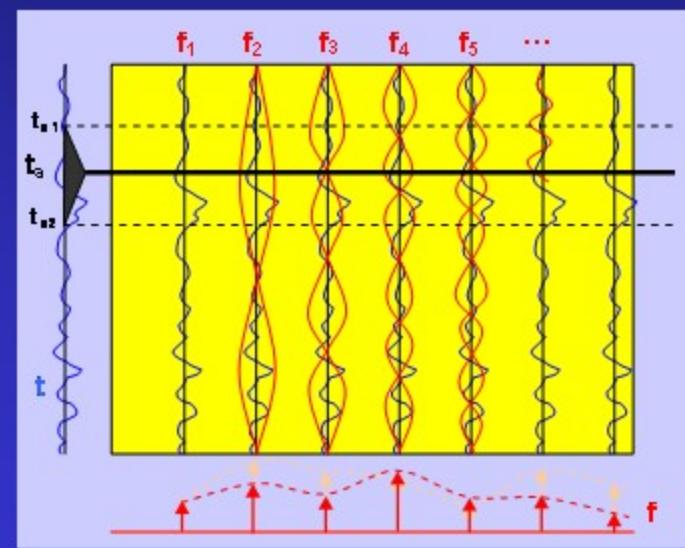
$$c(t) = \int_{t-t_1}^{t+t_1} (a\delta_{ta} + b\delta_{tb} + c\delta_{tc} + d\delta_{tc})\varsigma(\tau)d\tau = a\varsigma(t_a) + b\varsigma(t_b) + c\varsigma(t_c) + d\varsigma(t_c)$$

Fast-filtering, Gabor transform

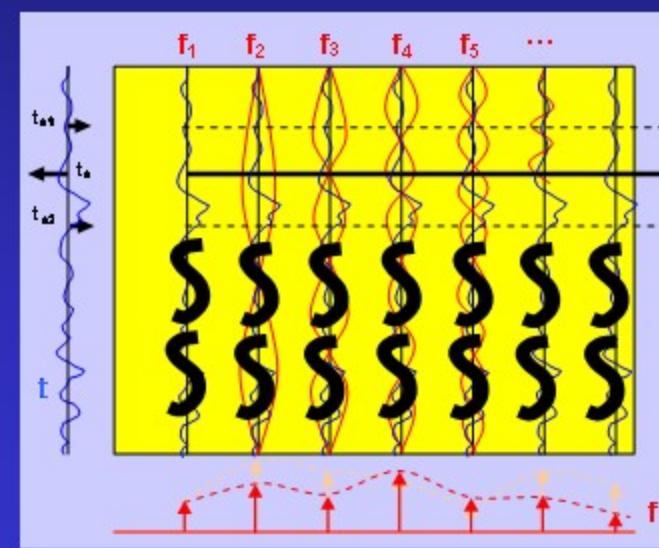
$$S(\omega, t) = \int_{t-t_1}^{t+t_1} w(\tau - t) g(\tau) e^{-i\omega\tau} d\tau$$



DFT



Spectrum one window



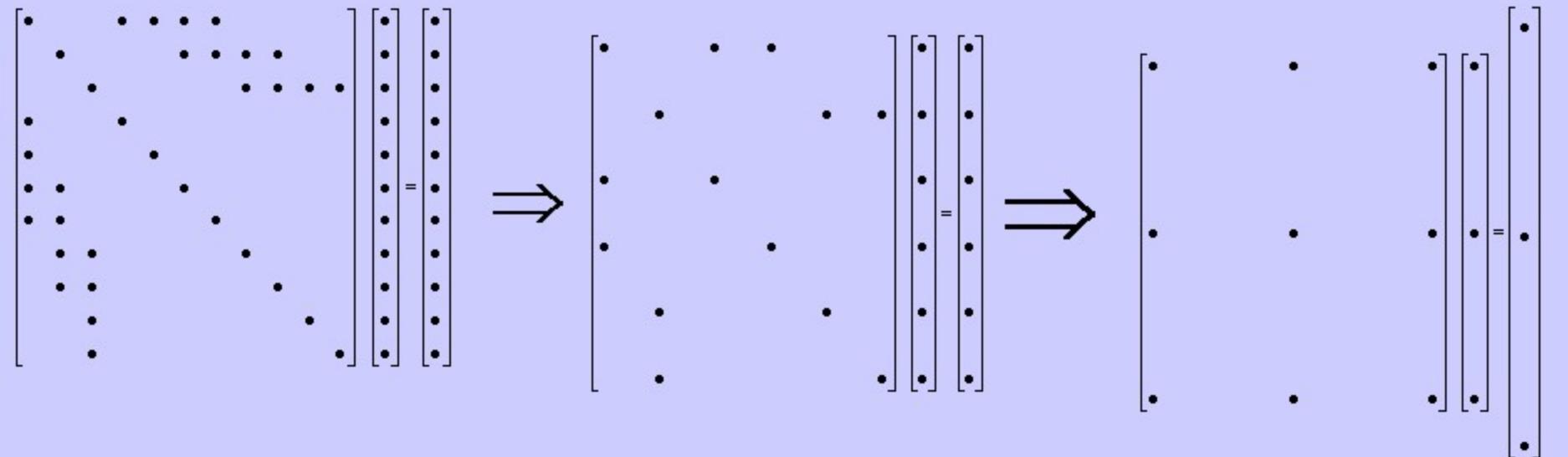
Fast-filter one window

Millar multigrid...Gauss-Seidel

Down sample



Solve the lower frequencies first



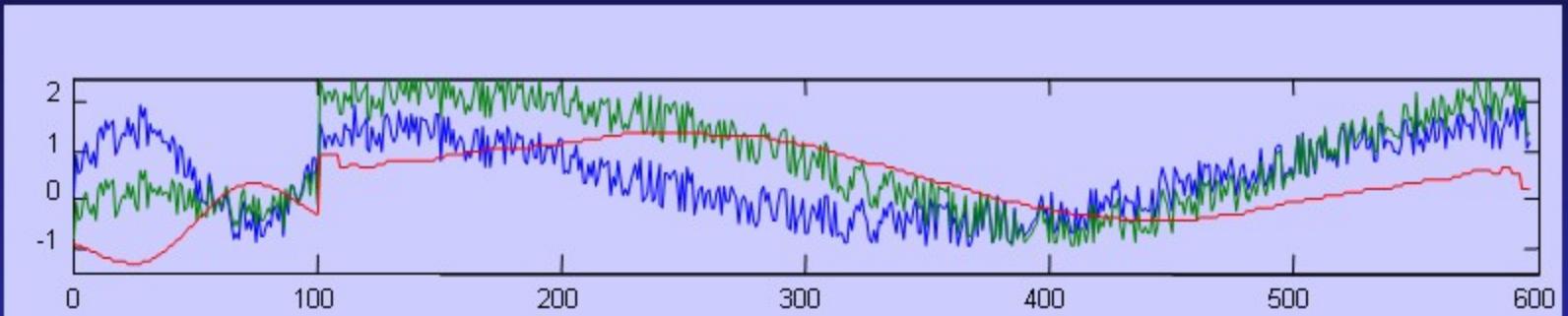
Add in the higher frequencies



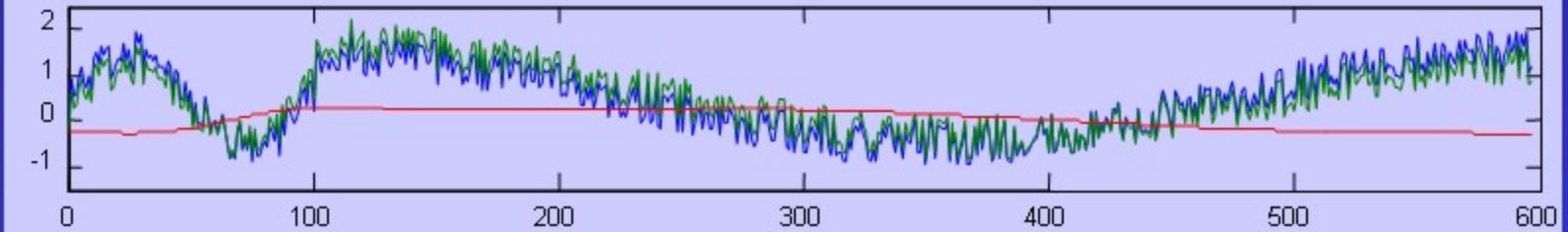
Interpolate

Surface consistent statics

Conventional GS

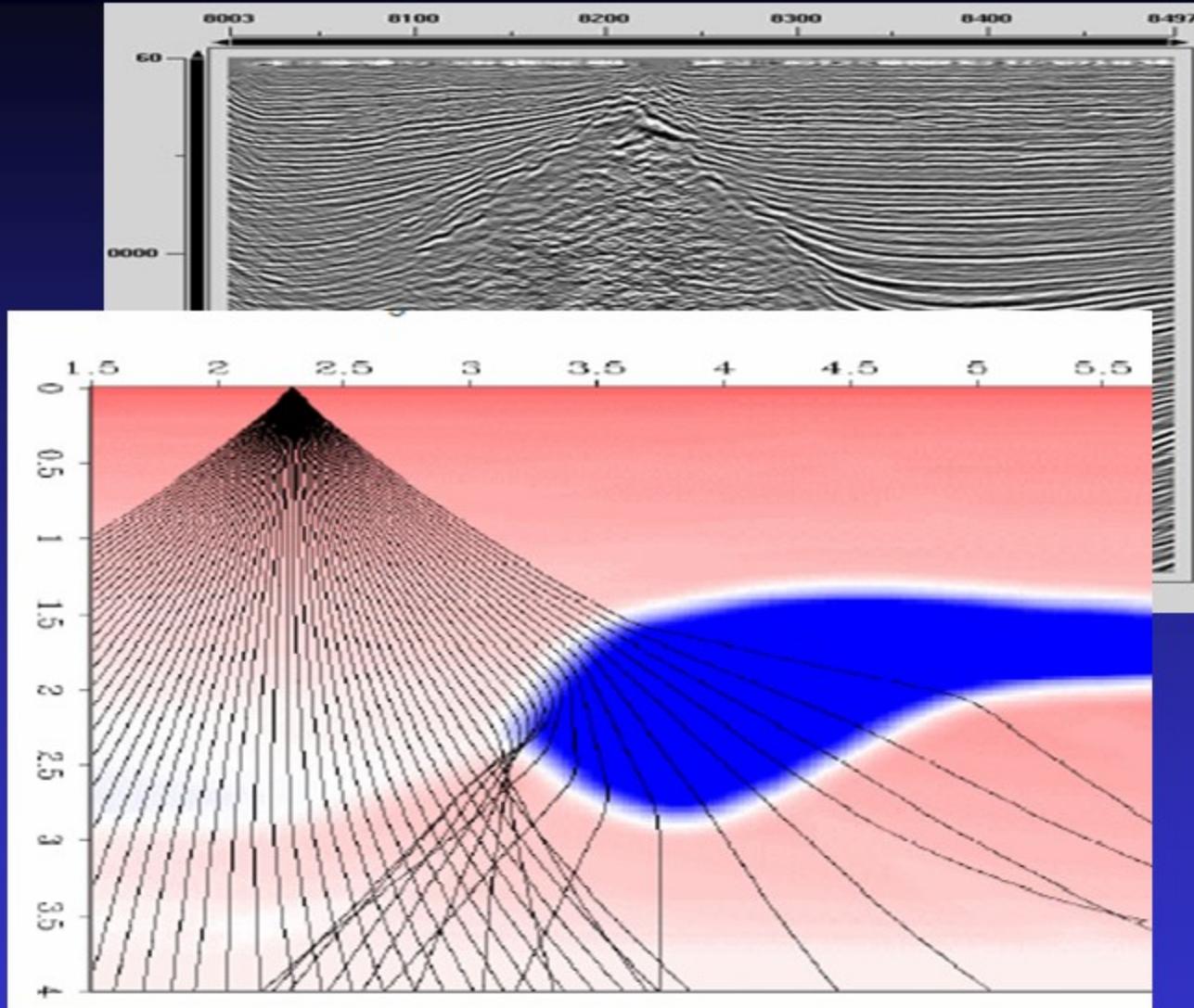


Multigrid



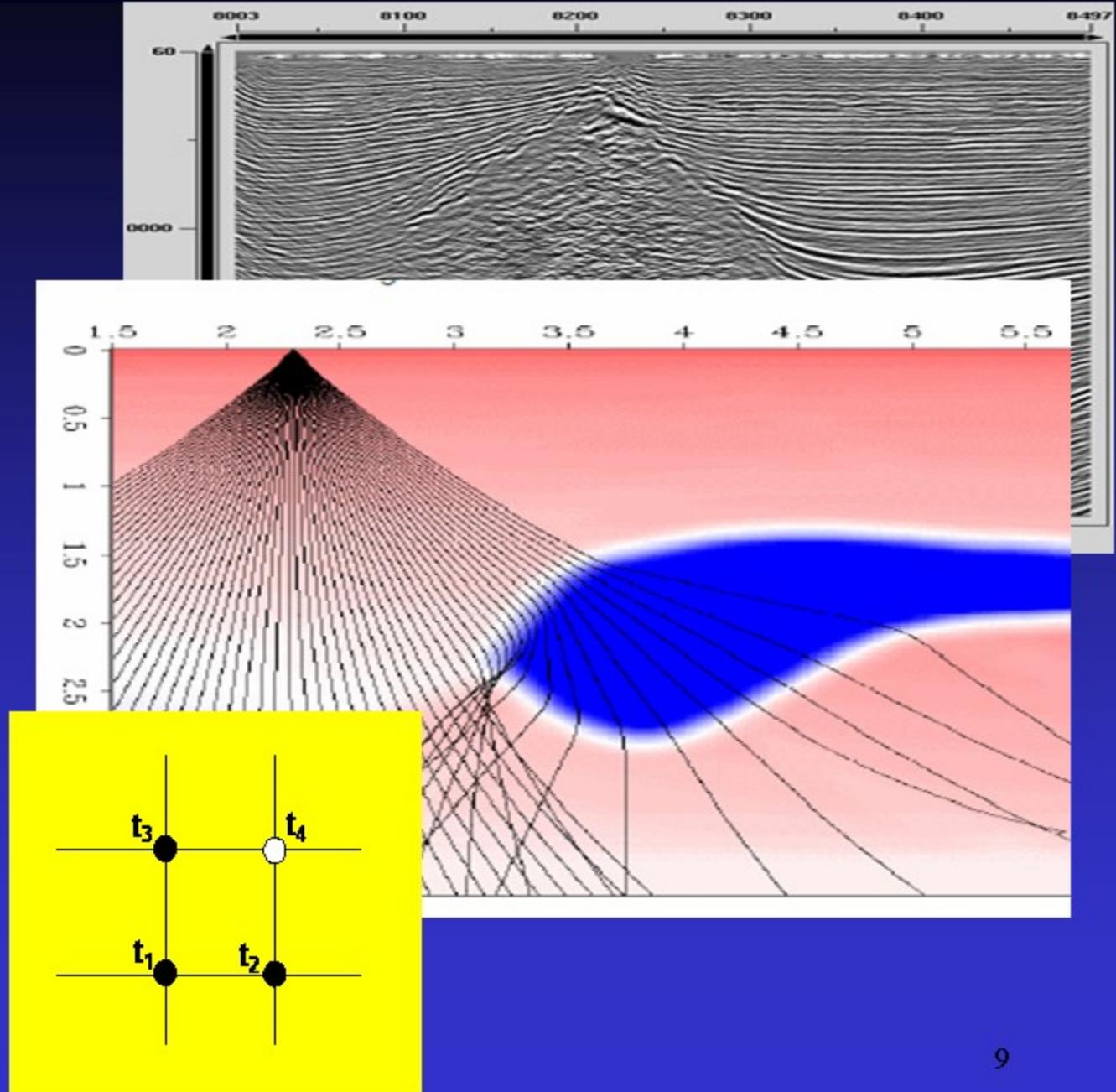
Synopsis

- Depth Migration
 - Kirchhoff migration
 - Ray tracing
 - Gridded traveltimes



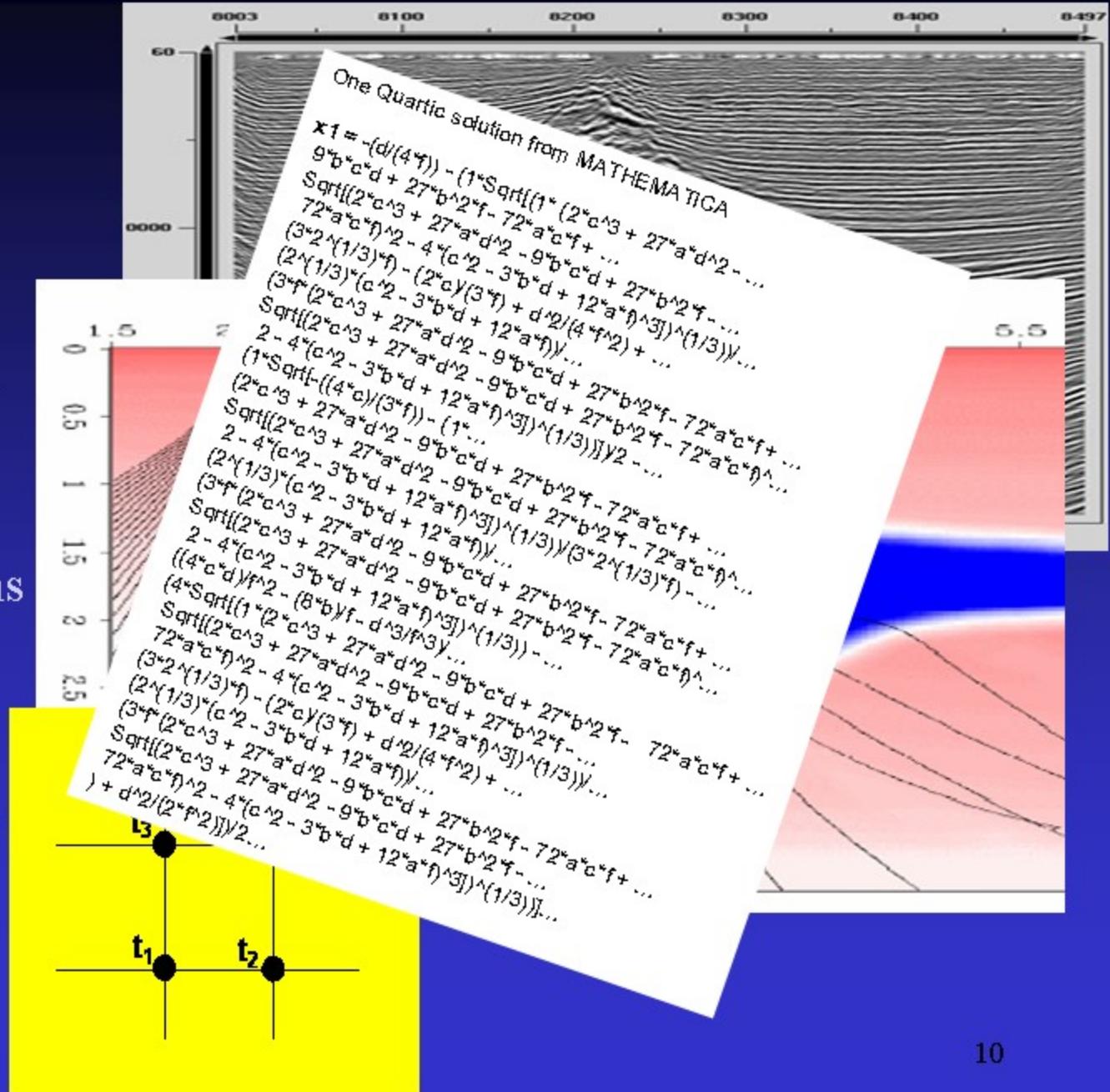
Synopsis

- Depth Migration
 - Kirchhoff migration
 - Ray tracing
 - Gridded traveltimes
- Three point Solutions
 - Plane wave
 - Vidale method
 - Circular assumption



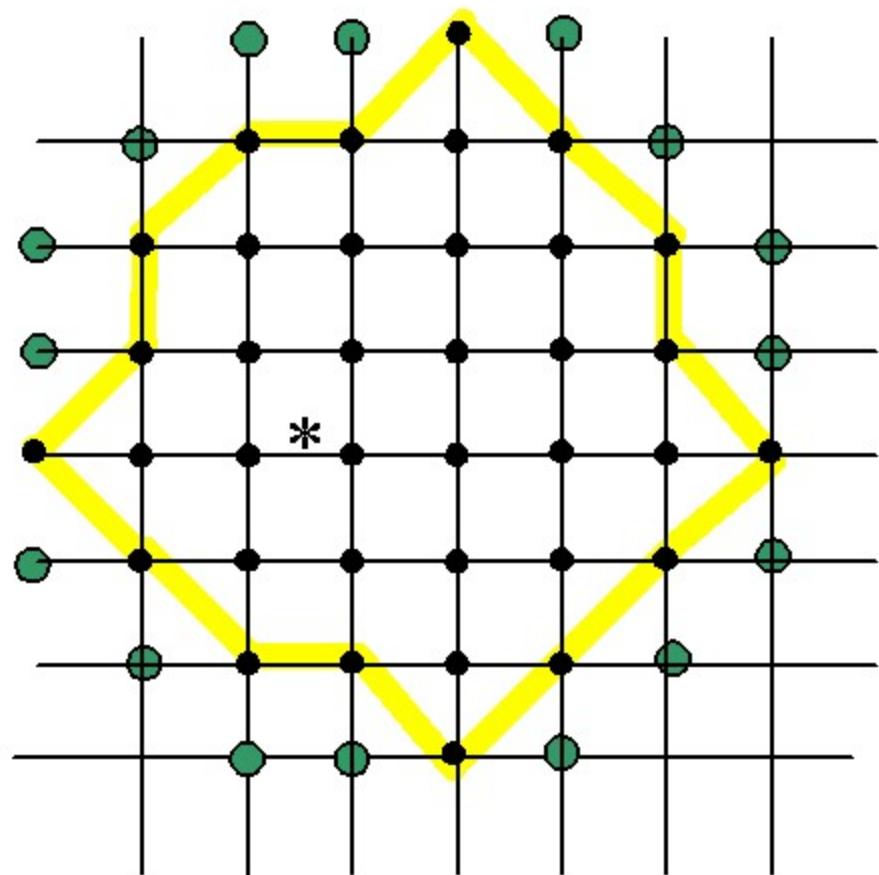
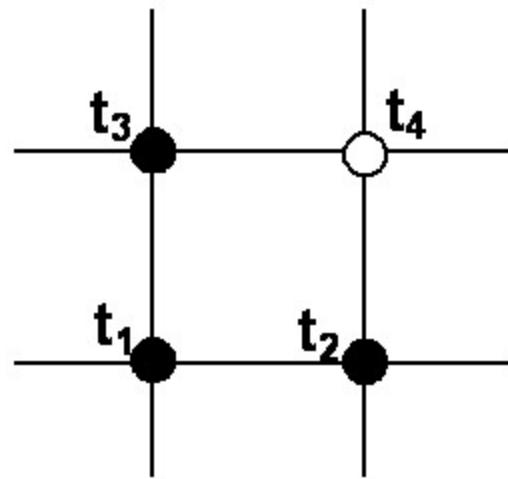
Synopsis

- Depth Migration
 - Kirchhoff migration
 - Ray tracing
 - Gridded traveltimes
- Three point Solutions
 - Plane wave
 - Vidale method
 - Circular assumption
- Computation
 - Quartic solution
 - Iterative solution



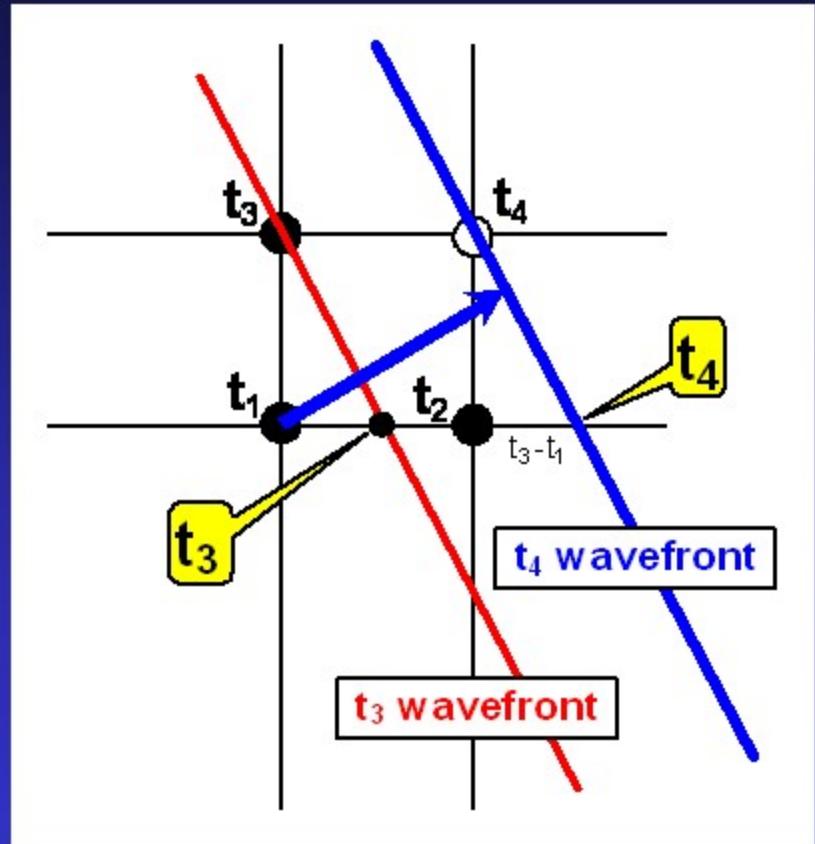
Update minimum times

**Wavefront tends to
circular shape**



Plane wave assumptions

$$t_4 = t_2 + t_3 - t_1$$



Vidale's method

Eikonal equation
Finite difference
solution of

$$\left(\frac{dt}{dx}\right)^2 + \left(\frac{dt}{dz}\right)^2 = \frac{1}{v^2}$$

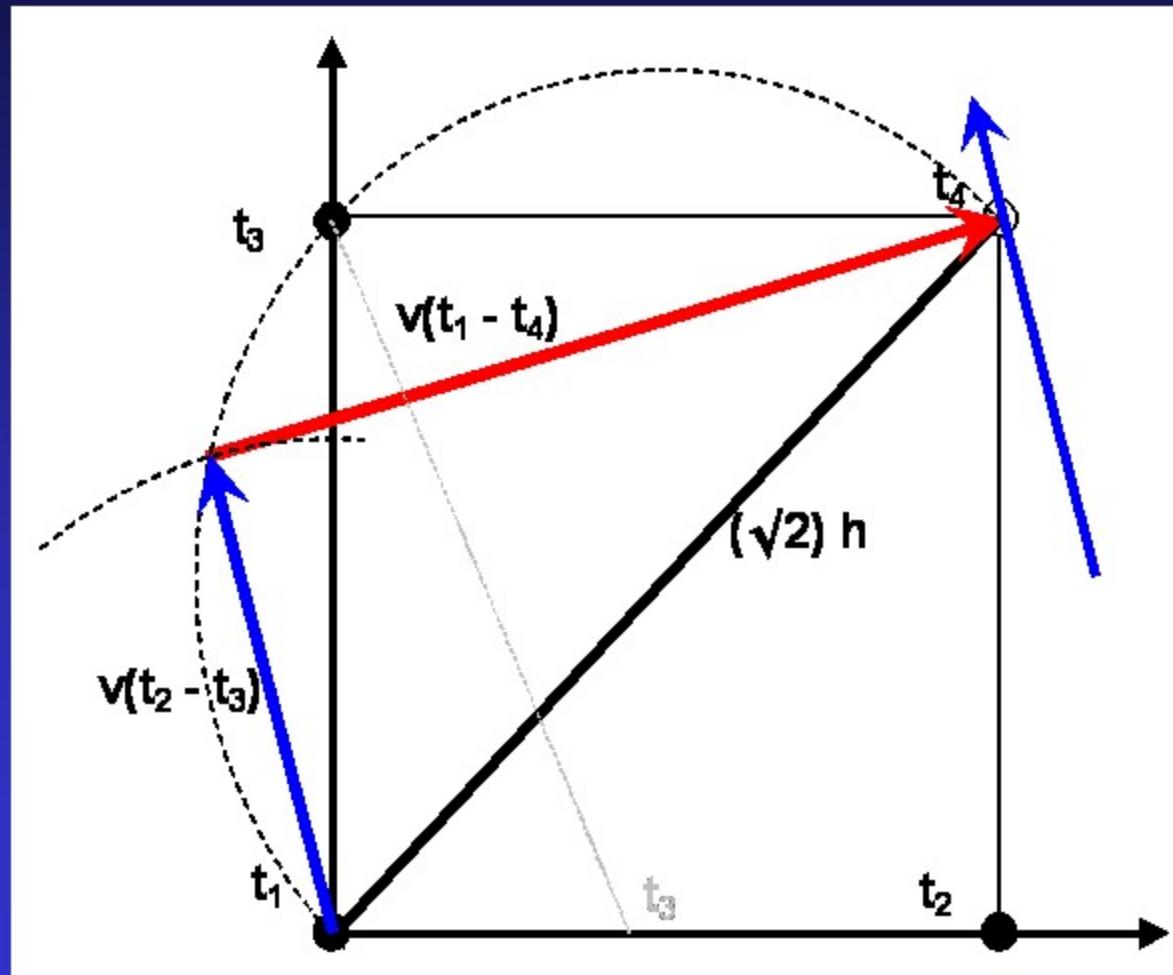
Basic solution in
common use

$$t_4 = t_1 + \sqrt{\frac{2h^2}{v^2} - (t_3 - t_2)^2}$$

Vidale's method

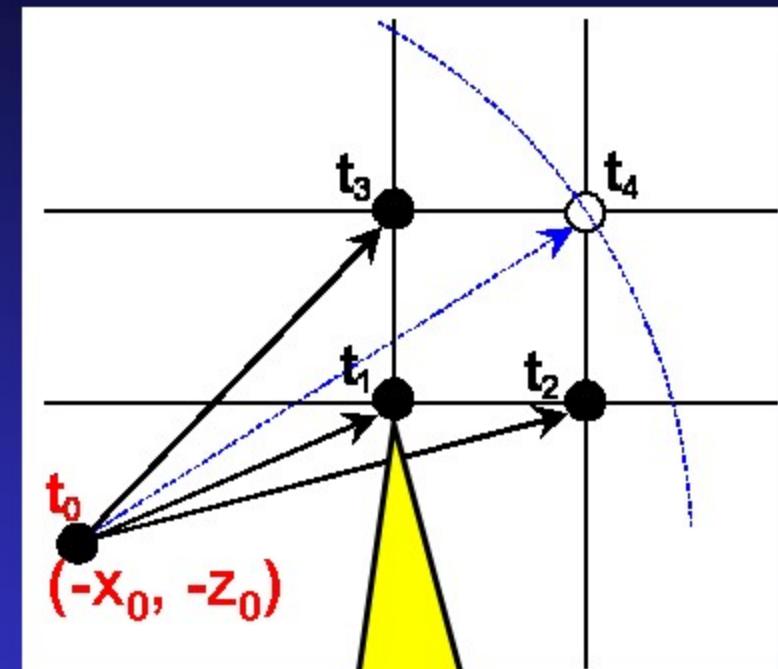
$$t_4 = t_1 + \sqrt{\frac{2h^2}{v^2} - (t_3 - t_2)^2}$$

$$(t_4 - t_1)^2 + (t_3 - t_2)^2 = \frac{2h^2}{v^2}$$



Curved wavefront assumption

1. Constant velocity of the square
2. Estimate center of curvature or apparent source location (x_0, z_0)
3. Define the time t_0
4. Get time t_4



(x, z) origin

The quartic equation

$$x^4 \left(\frac{b_1^4}{a_1^2} - \frac{2b_1^2 b_2^2}{a_1^2 a_2^2} + \frac{b_2^4}{a_2^4} \right)$$

$$+ x^3 \frac{4b_1^2}{a_1^2} \left(\frac{b_2^2}{a_2^2} - b_1^2 \right)$$

$$+ x^2 \left(+ \frac{6b_1^4}{a_1^4} + \frac{2b_2^4}{a_2^2} - \frac{2b_1^4}{a_1^2} - \frac{2b_1^2 b_2^2}{a_1^2} - \frac{2b_1^2 b_2^2}{a_1^2 a_2^2} + \frac{2b_1^2 b_2^2}{a_2^2} - \frac{2b_1^2}{a_1^2} - \frac{2b_2^2}{a_2^2} \right)$$

$$+ x \frac{4b_1^2}{a_1^2} \left(1 - \frac{b_1^2}{a_1^2} + b_1^2 + b_2^2 \right)$$

$$+ \left(b_1^4 + b_2^4 + 2b_1^2 - 2b_2^2 + 2b_1^2 b_2^2 + \frac{b_1^4}{a_1^4} - \frac{2b_1^4}{a_1^2} - \frac{2b_1^2}{a_1^2} - \frac{2b_1^2 b_2^2}{a_1^2} + 1 \right)$$

$$= 0$$

$$a + xb + x^2c + x^3d + x^4f = 0$$

A quartic solution

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File Edit View Insert Format Tools Table MathType Window Help

Type a question for help

Final Showing Markup Show

One Quartic solution:

```
(d/(4*f))^(1/4)*Sqrt[(1*(c^3+27*a*d^2)^2-9*b*c*d+27*b^2*f)^(1/3)]*Sqrt[(2*c^3+27*a*d^2-9*b*c*d+27*b^2*f)^(1/3)]^(1/2)*(c^2-3*b*d+12*a*f)^3]^(1/3))/(-3*2^(1/3)*(c^2-3*b*d+12*a*f))^(1/2)*((2*c/(3*f)+d^2/(4*f^2))+(-2^(1/3)*(c^2-3*b*d+12*a*f))^(1/2)*(2*c^3+27*a*d^2-9*b*c*d+27*b^2*f-72*a*c*f)^(1/3)*Sqrt[(2*c^3+27*a*d^2-9*b*c*d+27*b^2*f-72*a*c*f)^(1/3)]^(1/2)*(-2^(1/3)*(c^2-3*b*d+12*a*f)^3)^(1/3))/2^(1/2)*(1*Sqrt[-(4*c)/(3*f)]^(1/4)*(2*c^3+27*a*d^2-9*b*c*d+27*b^2*f-72*a*c*f)^(1/3)*Sqrt[(2*c^3+27*a*d^2-9*b*c*d+27*b^2*f-72*a*c*f)^(1/3)]^(1/2)*(-2^(1/3)*(c^2-3*b*d+12*a*f)^3)^(1/3))/(3*2^(1/3)*f)^(1/2)*(2^(1/3)*(c^2-3*b*d+12*a*f))^(1/2)*(3*f^2/(2^(1/2)*(B*b)/f-d^3/f^3))-4*Sqrt[(1*(2*c^3+27*a*d^2-9*b*c*d+27*b^2*f)^2-72*a*c*f)^(1/3)]*Sqrt[(2*c^3+27*a*d^2-9*b*c*d+27*b^2*f)^(1/3)]^(1/2)*(-72*a*c*f)^(1/2)*(-4*(c^2-3*b*d+12*a*f)^3)^(1/3)/(-4*c*d)/f^2*(B*b)/f-d^3/f^3)-4*Sqrt[(1*(2*c^3+27*a*d^2-9*b*c*d+27*b^2*f)^2-72*a*c*f)^(1/3)]*Sqrt[(2*c^3+27*a*d^2-9*b*c*d+27*b^2*f)^(1/3)]^(1/2)*(3*2^(1/3)*f)^(1/2)*(2*c/(3*f)+d^2/(4*f^2))+(-2^(1/3)*(c^2-3*b*d+12*a*f))^(1/2)*(3*f^2/(2^(1/2)*(2*f^2)))^(1/2)
```

Formed from MATHEMATICA

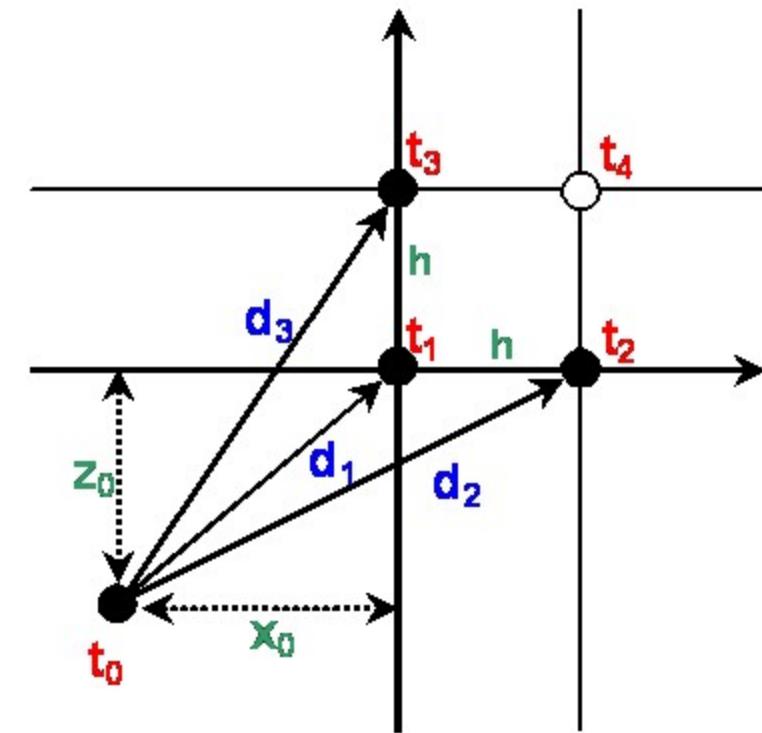
1.→Cell→Convert To→Input Form, then

2.→Edit→Copy As→Text

Page 10 Sec 1 10/11 At: 2.5" Ln 7 Col 36 REC TRK EXIT OWR English (U.S)

17

$$\left\{ \begin{array}{l} t_0 = t_1 - \frac{1}{v} \sqrt{x_0^2 + z_0^2} \\ \\ x_0 = \frac{v^2}{2h} [2t_0(t_2 - t_1) - t_2^2 + t_1^2] + \frac{h}{2} \\ \\ z_0 = \frac{v^2}{2h} [2t_0(t_3 - t_1) - t_3^2 + t_1^2] + \frac{h}{2} \end{array} \right.$$



$$t_0 = t_1 - \frac{1}{v} \sqrt{\left\{ \frac{v^2}{2h} [2t_0(t_2 - t_1) - t_2^2 + t_1^2] + \frac{h}{2} \right\}^2 + \left\{ \frac{v^2}{2h} [2t_0(t_3 - t_1) - t_3^2 + t_1^2] + \frac{h}{2} \right\}^2}$$

Only variable is t_0 ...implicit

$$t_0 = t_1 - \frac{1}{v} \sqrt{\left\{ \frac{v^2}{2h} \left[2t_0(t_2 - t_1) - t_2^2 + t_1^2 \right] + \frac{h}{2} \right\}^2 + \left\{ \frac{v^2}{2h} \left[2t_0(t_3 - t_1) - t_3^2 + t_1^2 \right] + \frac{h}{2} \right\}^2}$$

Enter ... Newton Raphson

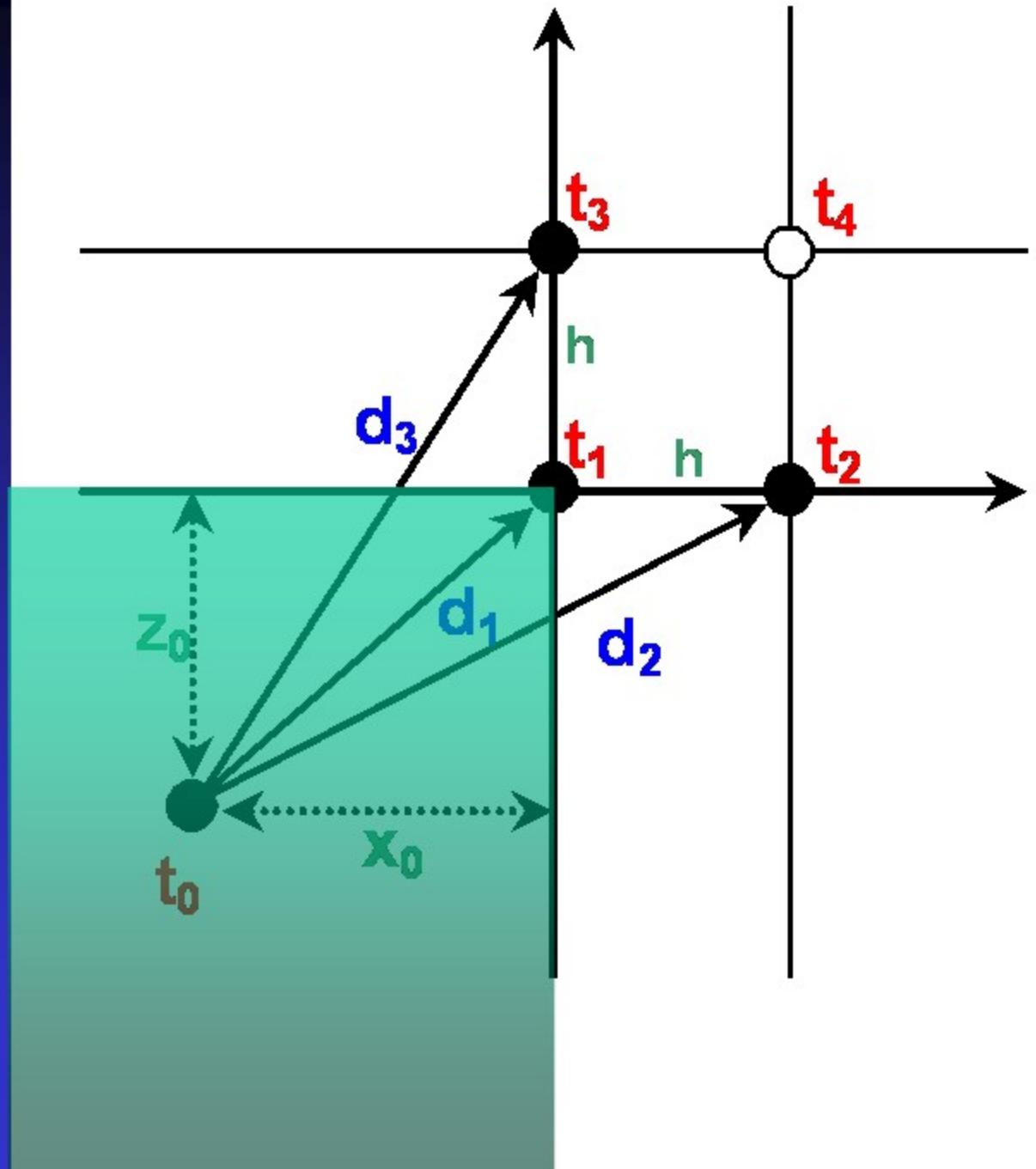
Find t_0 $f(t_0) \Rightarrow 0$

$$\frac{df(t_0)}{dt_0} = \sim$$

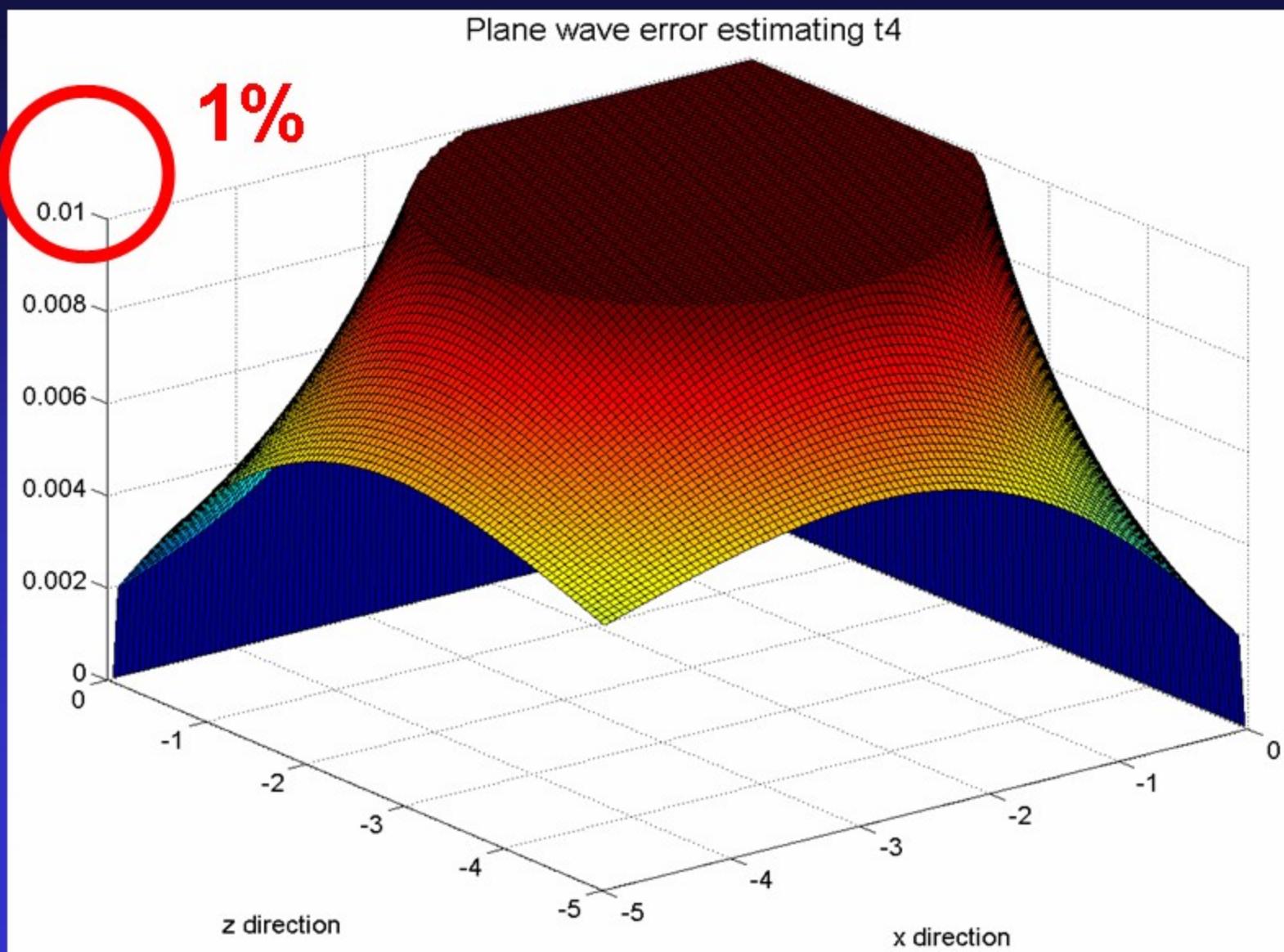
**Solution in
third
quadrant
(expanding
wavefront)**

$X = -5 \text{ to } 0$

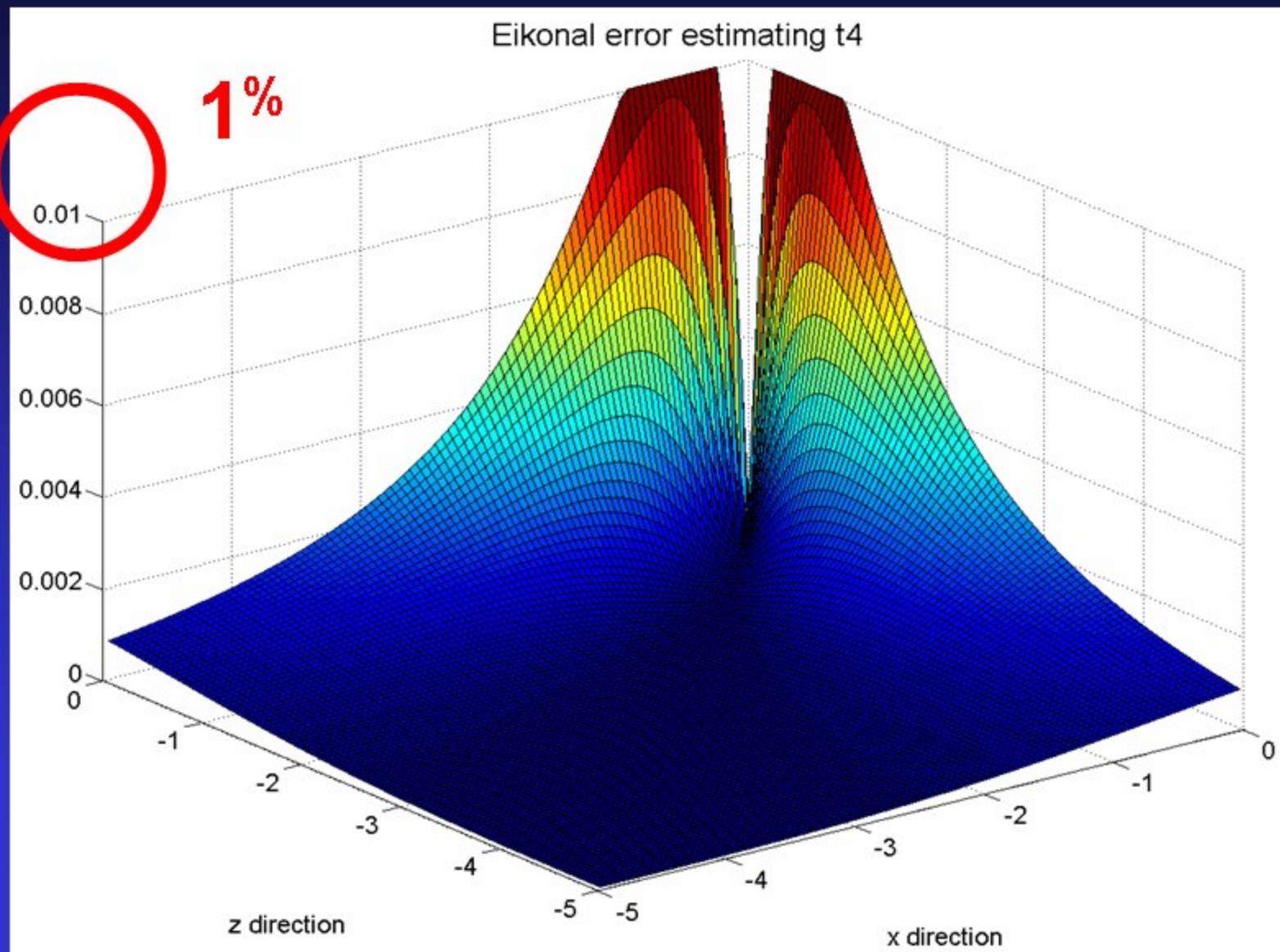
$Z = -5 \text{ to } 0$



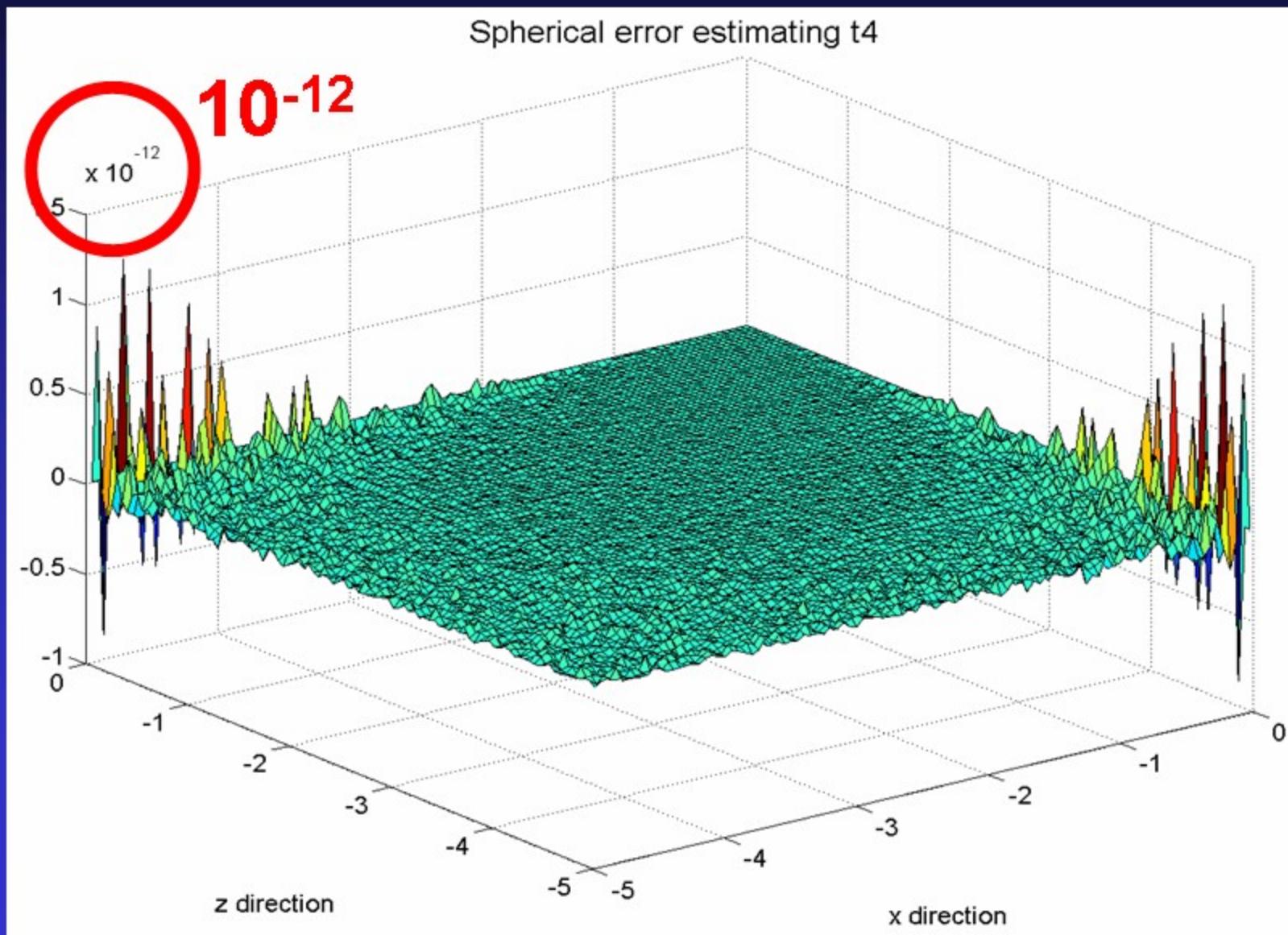
Plane wave solution



Vidale's solution

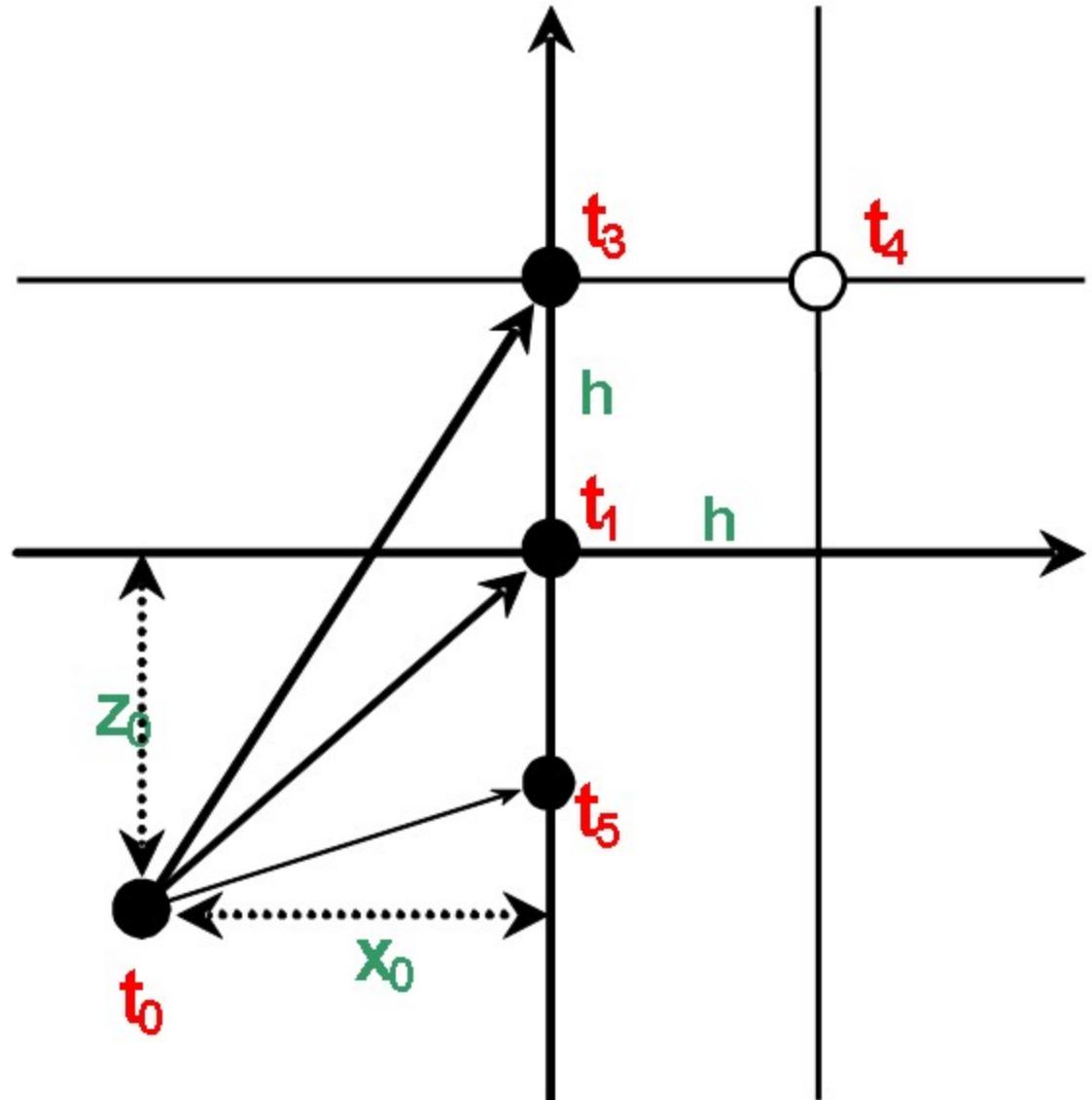


Iterative solution



Alternate solution

t_0, z_0, x_0

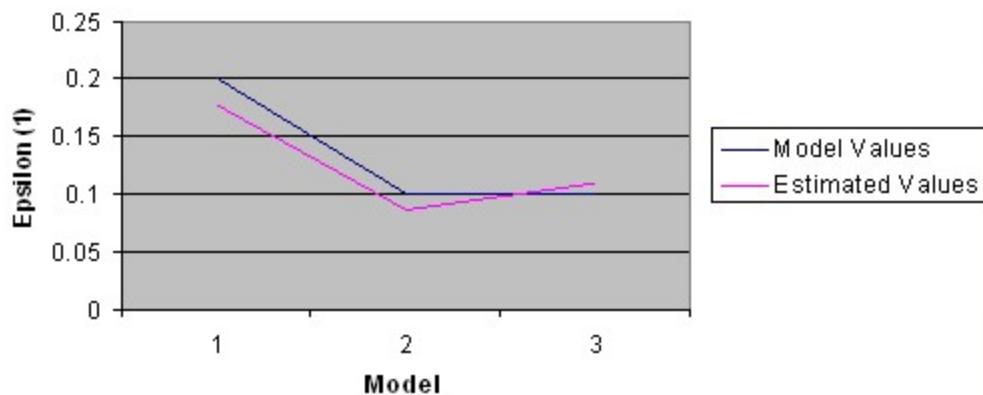


Pavan Elapavuluri: Estimation of anisotropy parameters

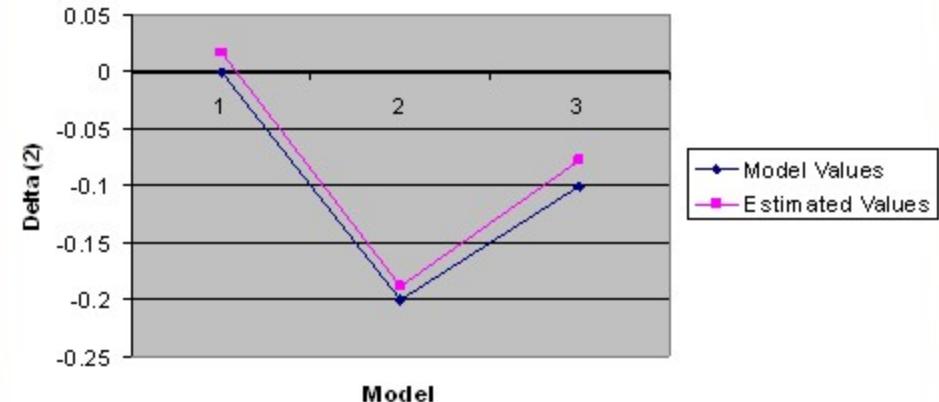
Thanks for the help Pat

Orthorhombic media

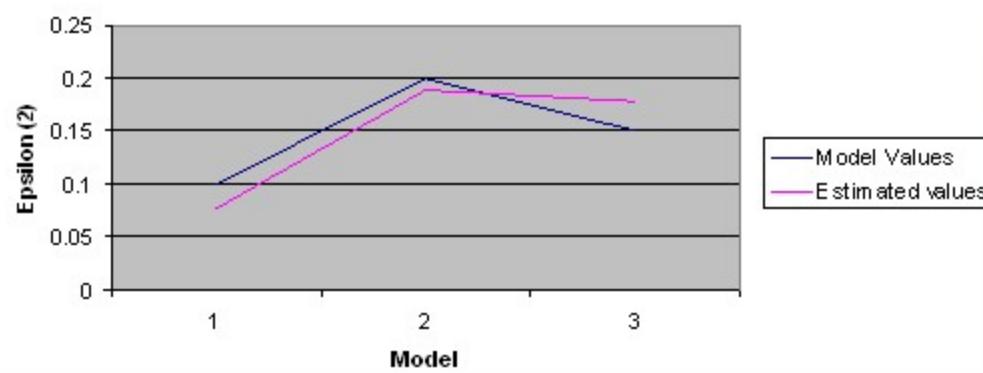
Comparision of estimated Epsilon (1) values



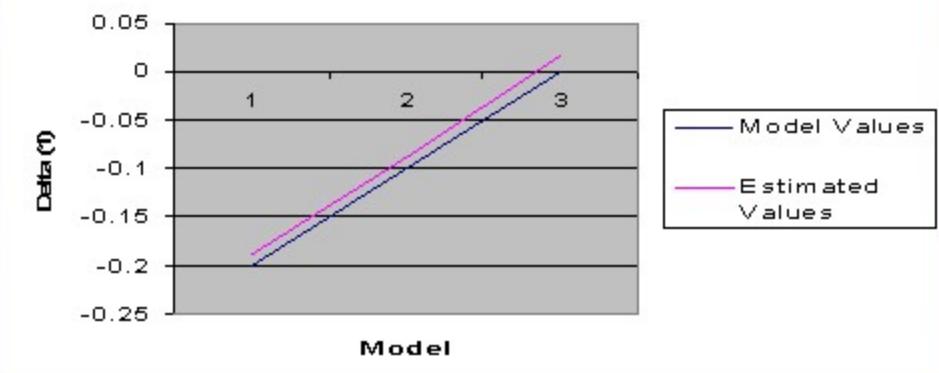
Comparison of estimated Delta (2) values



Comparision of estimated Epsilon(2) values



Comparison of estimated Delta (1) Values

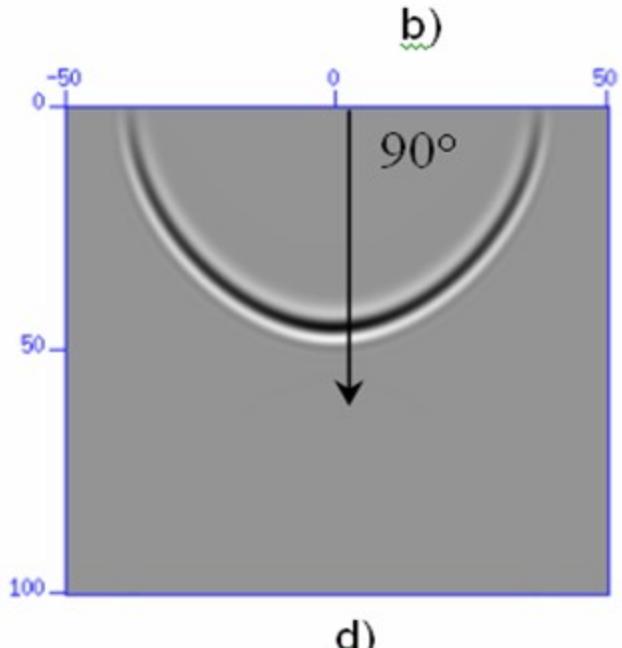
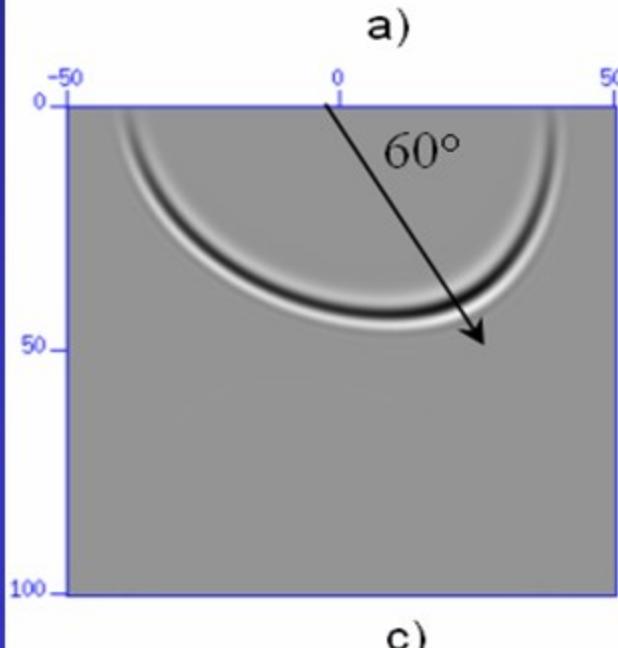
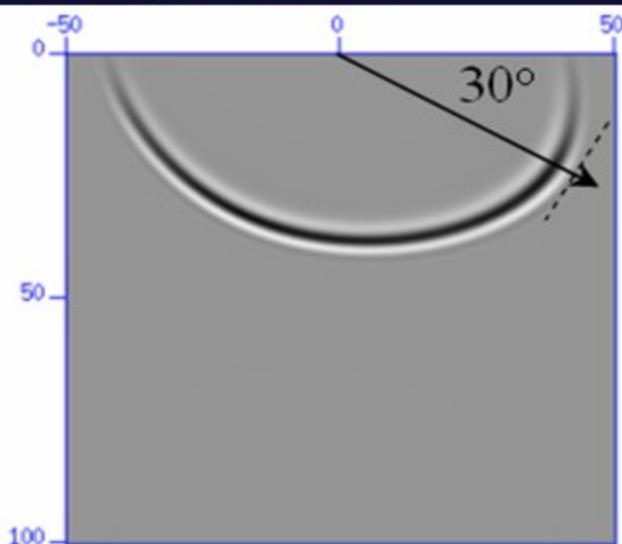
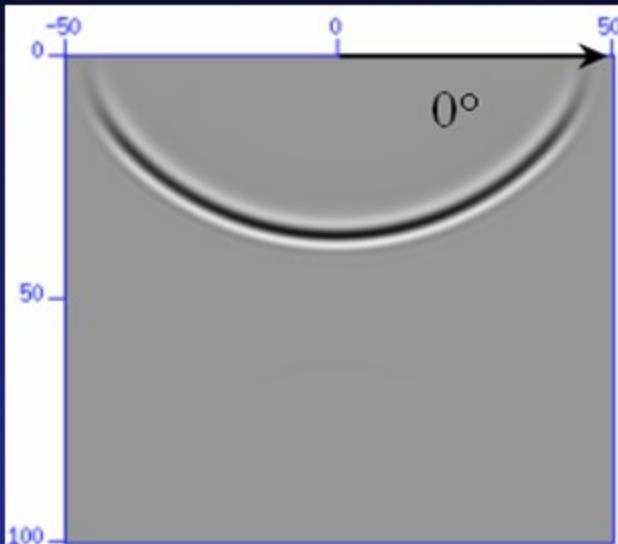


Anisotropic Reverse-Time and PSPI migration for Tilted TI Media

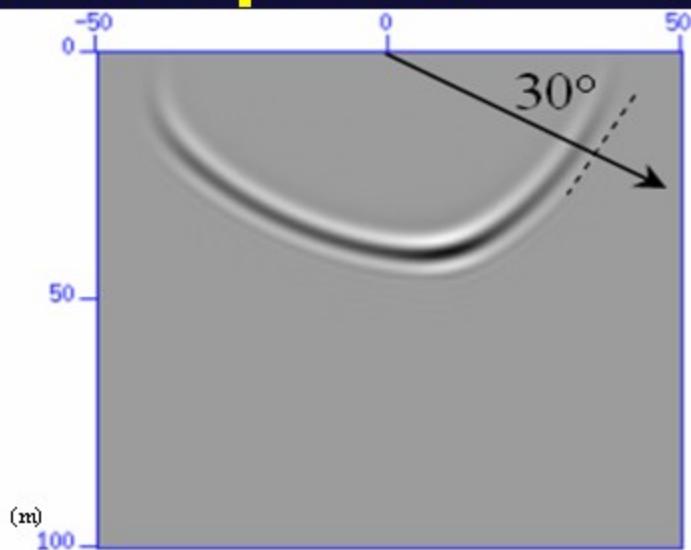
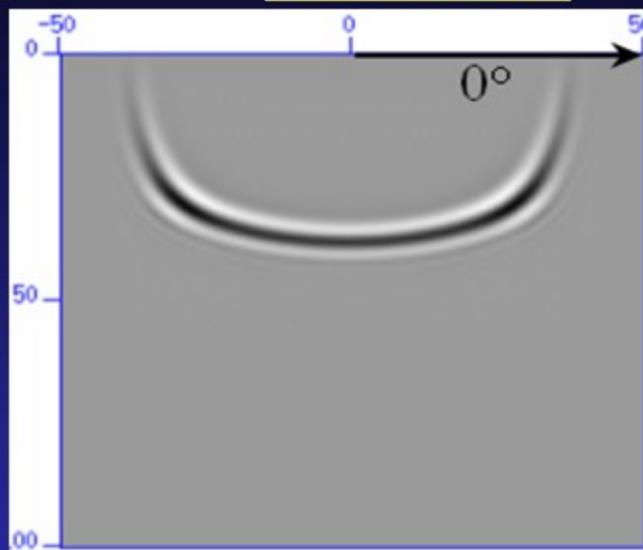
Xiang Du, John C. Bancroft,
Larry R. Lines and Don Lawton

Numerical example: RT migration (x , z)

P wave impulse response



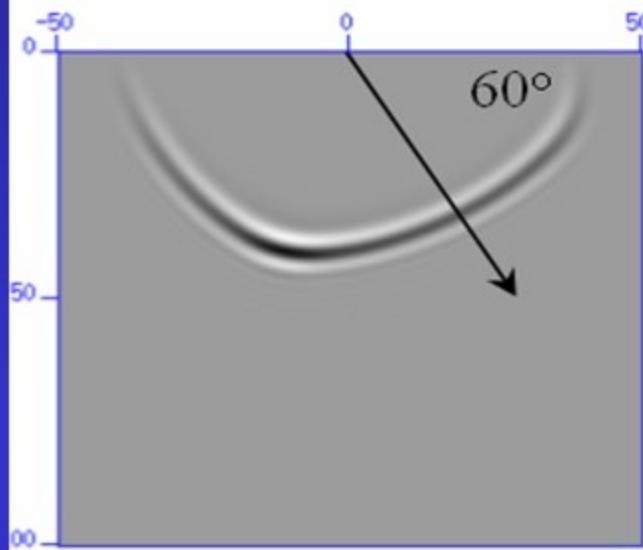
Numerical example: migration (x , z) S wave impulse response



a)

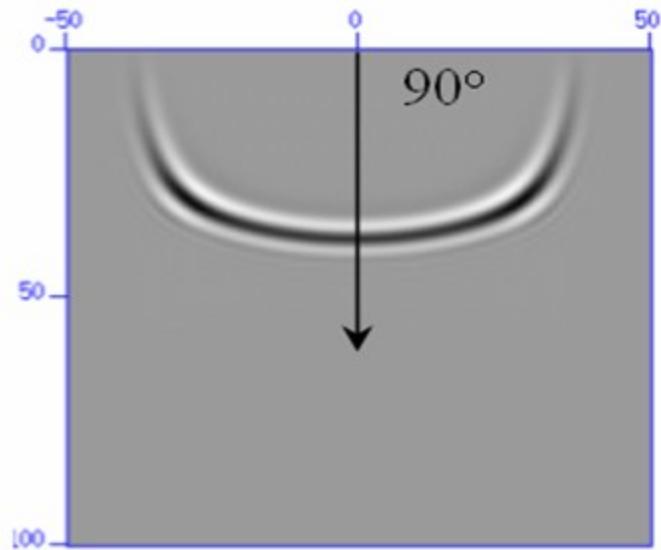
b)

60°

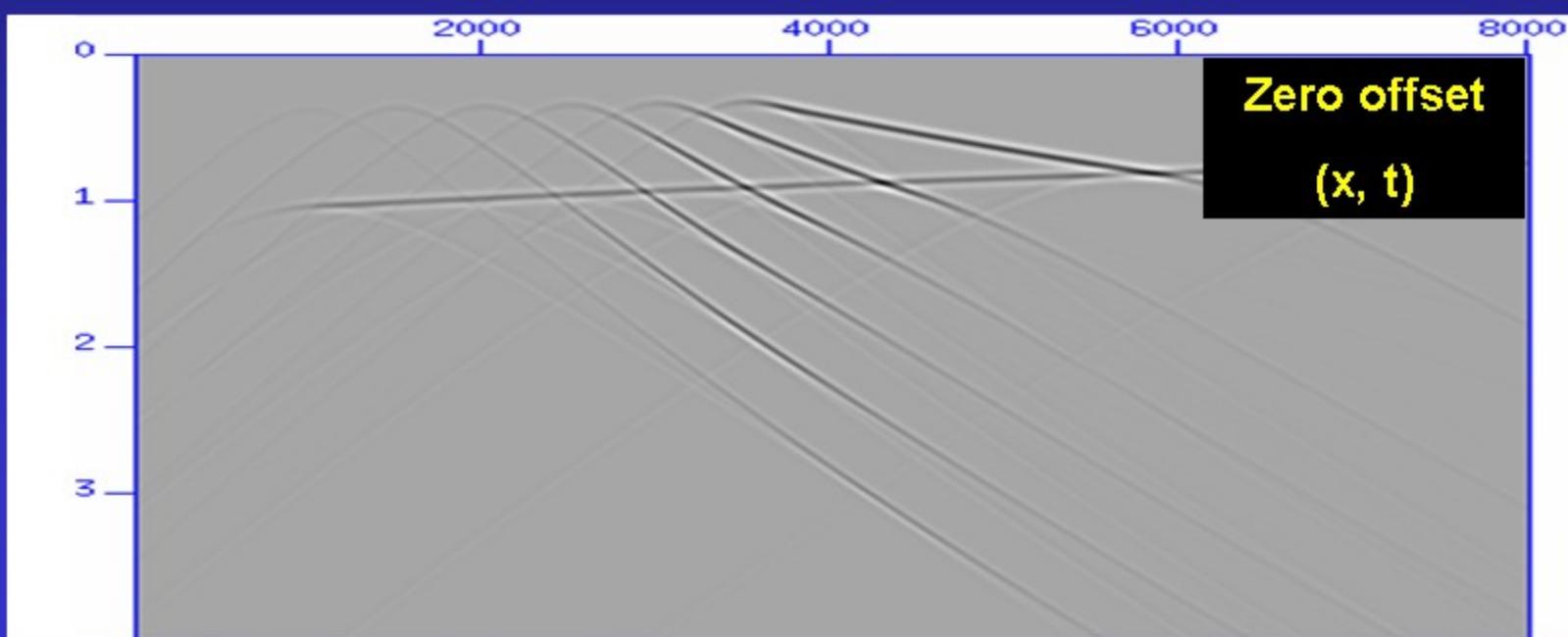
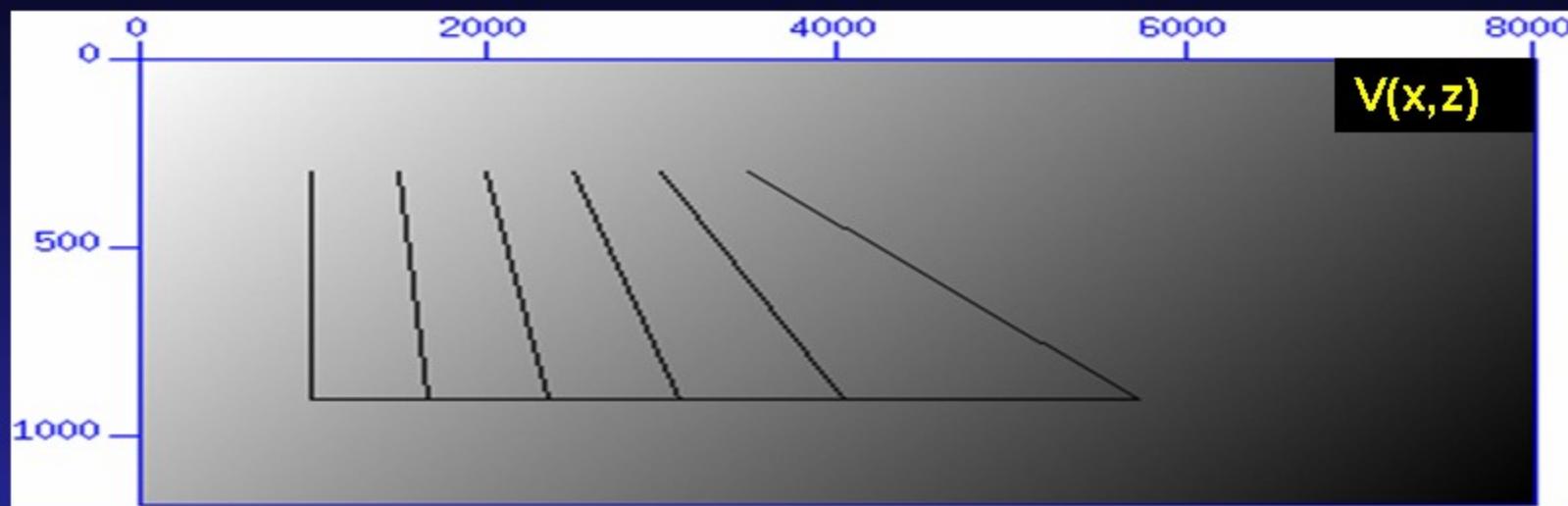


c)

d)

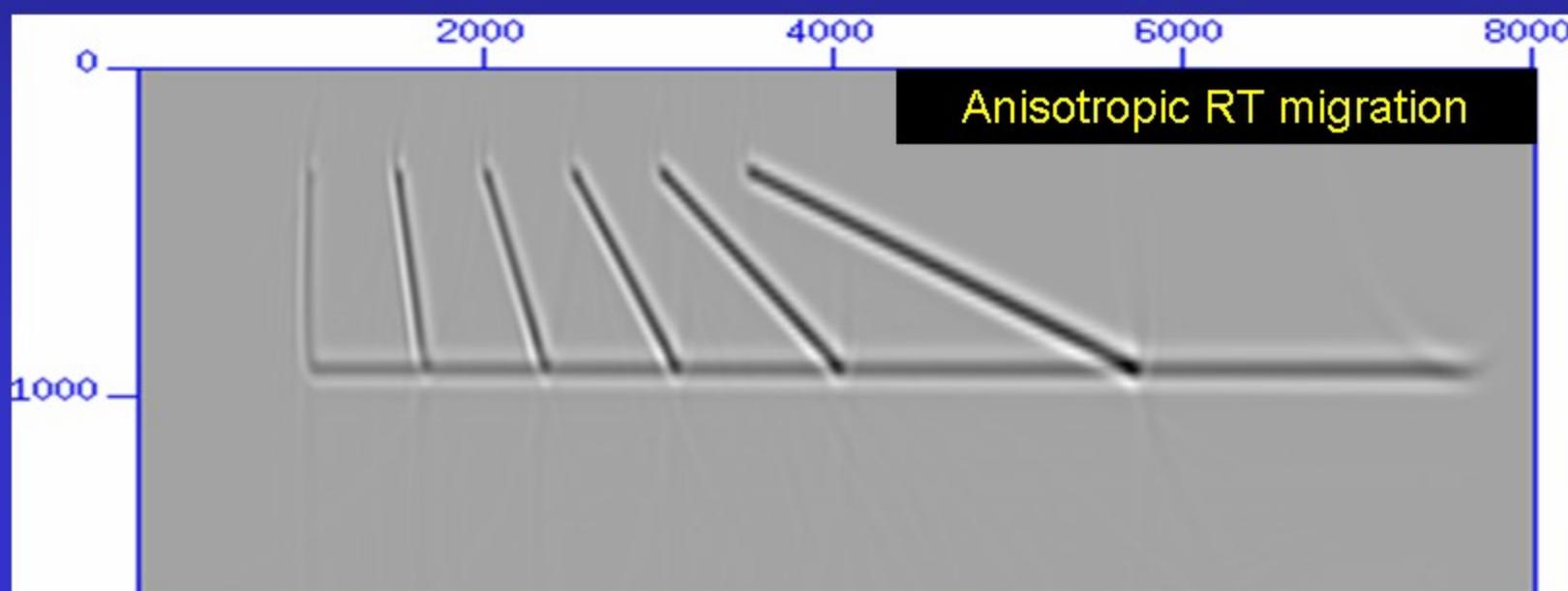
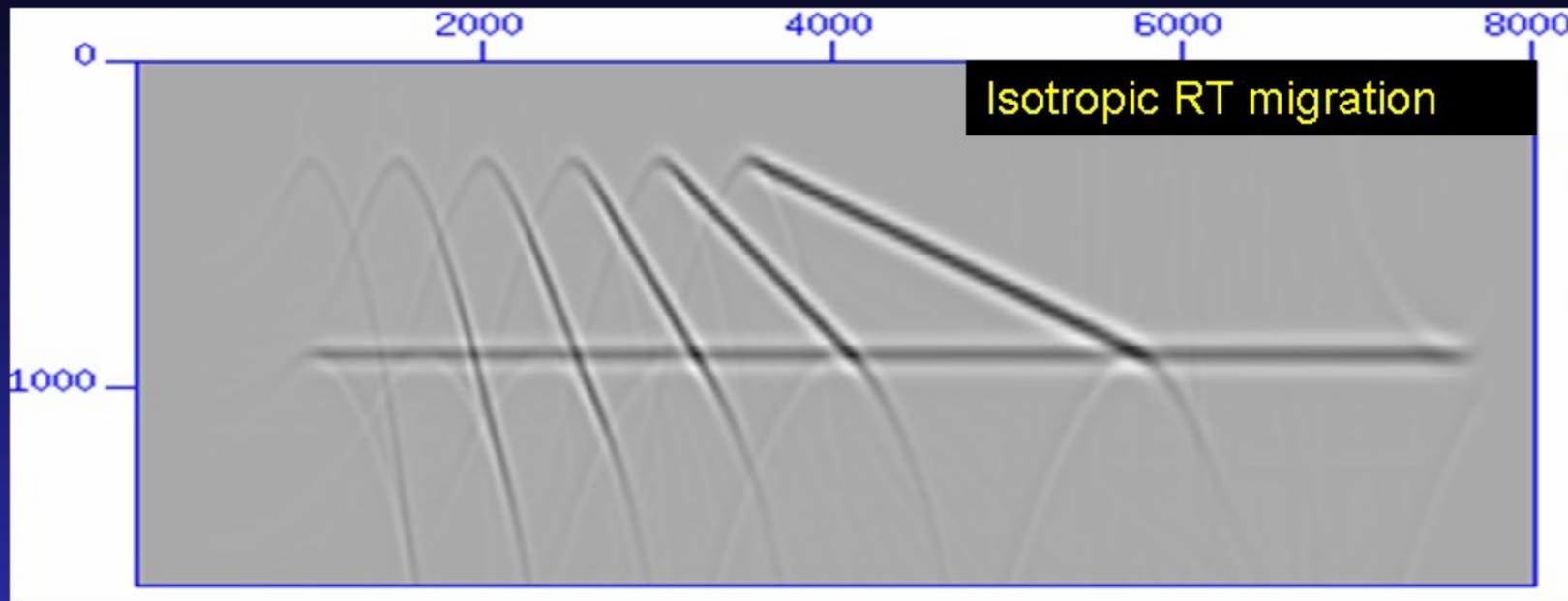


Numerical example: Variable A-velocity model

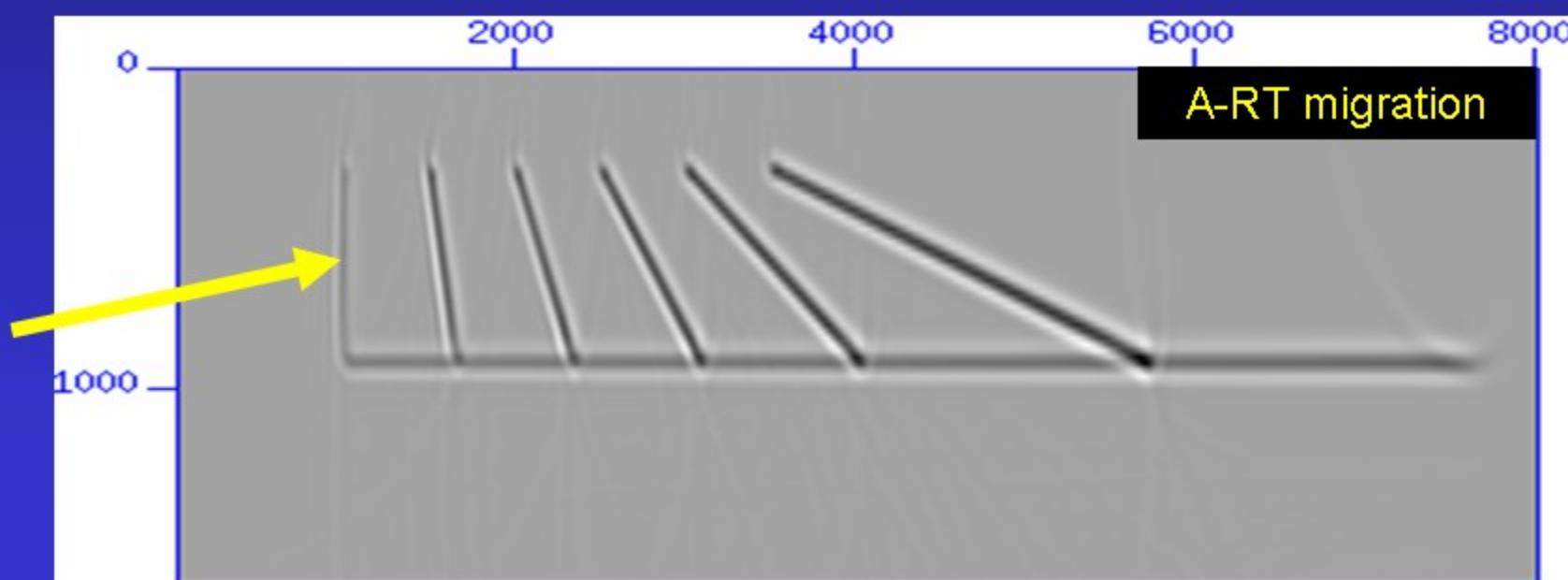
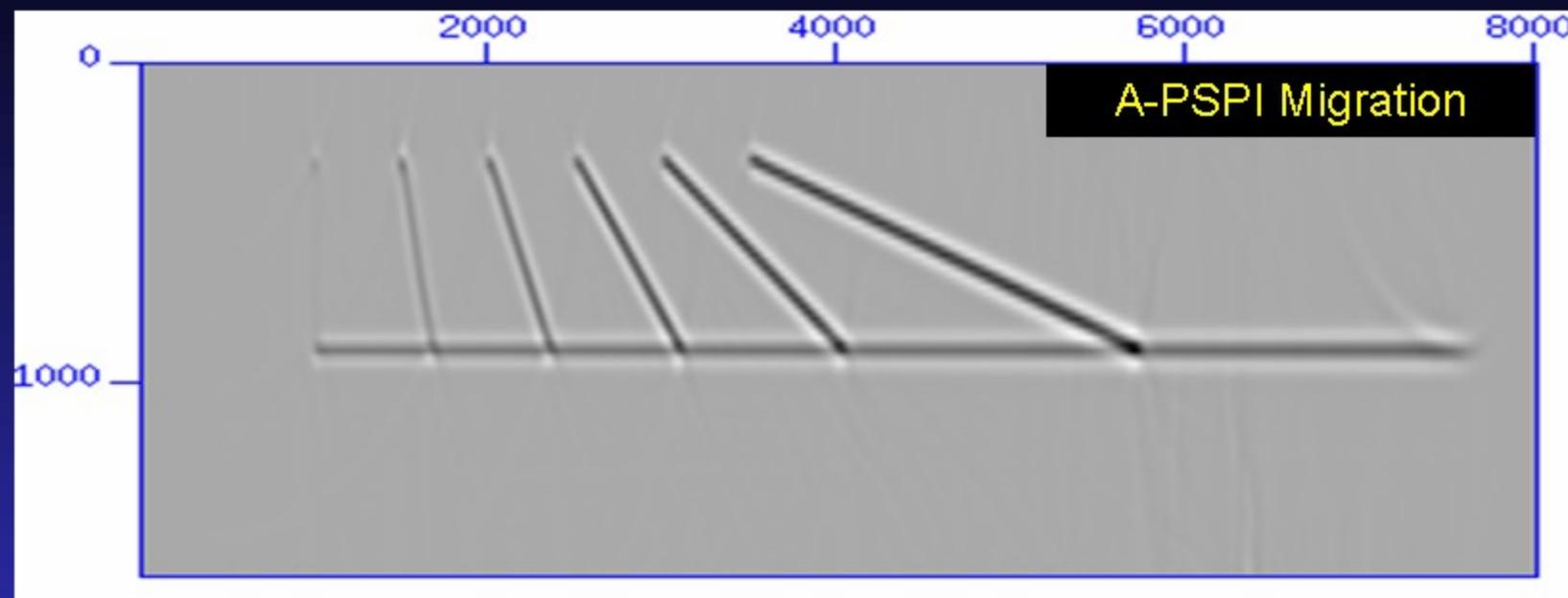


Synthetic zero-offset seismogram obtained using an SU code from Center for Wave Phenomena (CWP).

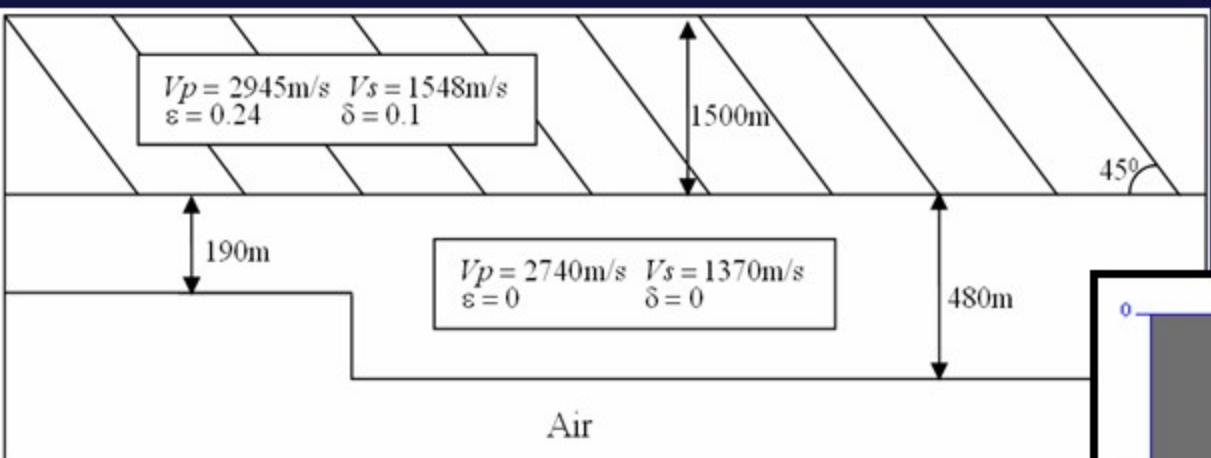
Comparison: Iso. and Aniso. Reverse time mig.



Comparison: A-PSPI A-RT

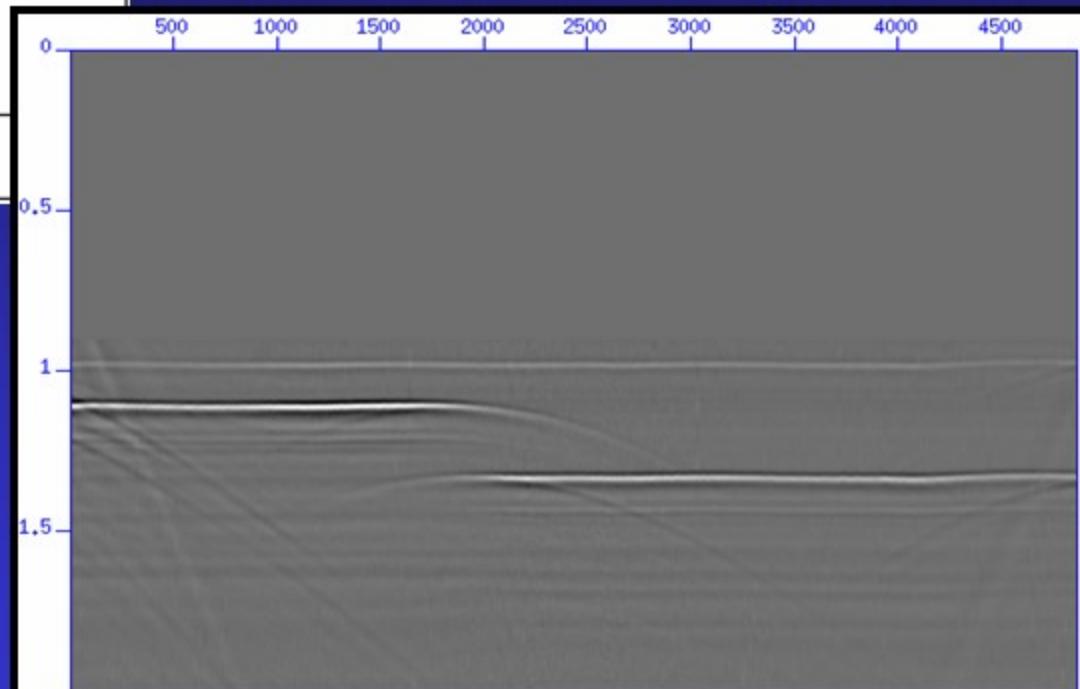


Physical model: reef model with a TTI overburden



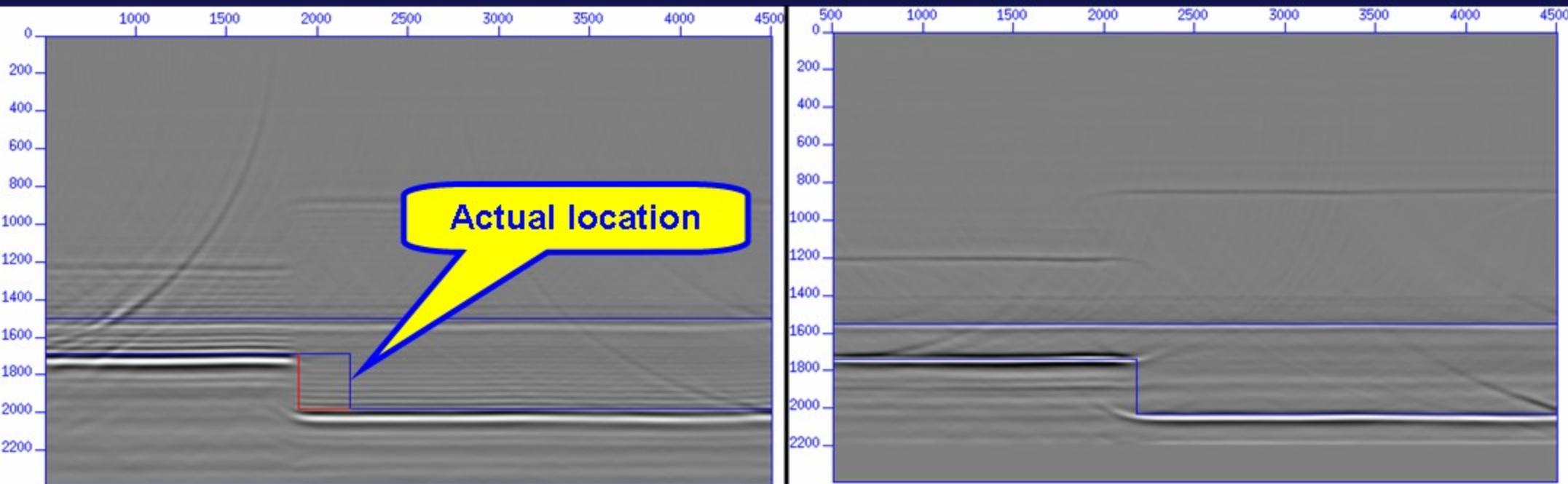
Isotropic “reef” with a TTI overburden

Dr. Don Lawton, Dr. Helen Isaac (models)



Zero-offset seismic section of reef model.

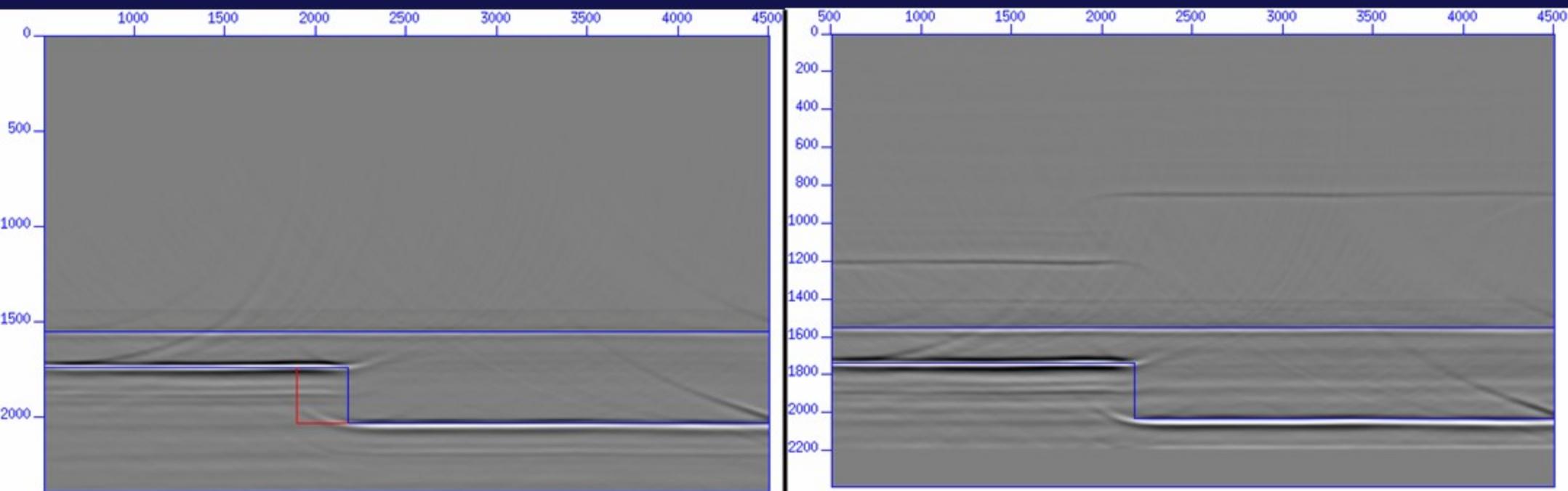
Physical model: Poststack migrations



Isotropic RT migration result of the 6th order accuracy for reef model.

Anisotropic RT migration.

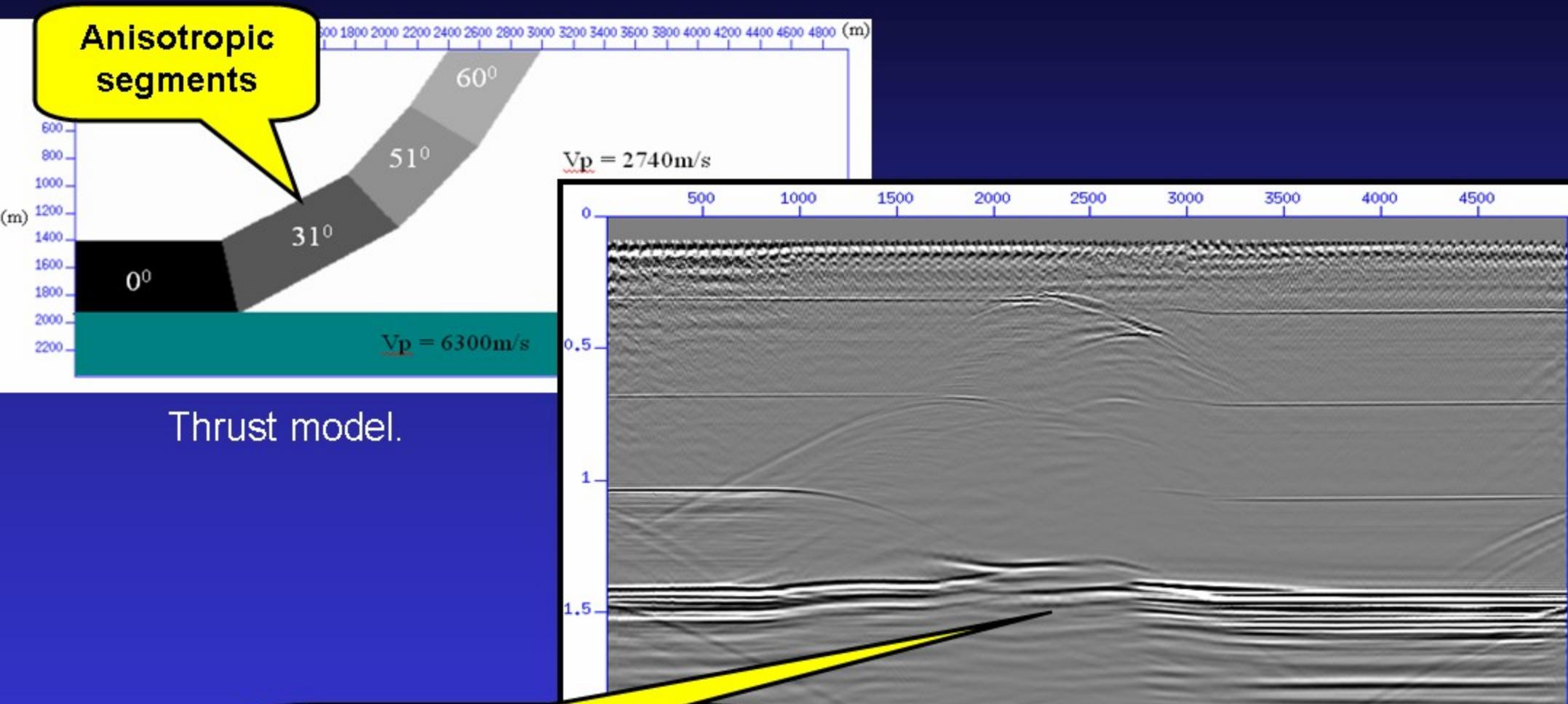
Physical model: Poststack migrations



Anisotropic **PSPI** migration.

Anisotropic **RT** migration.

Physical model: Thrust fault

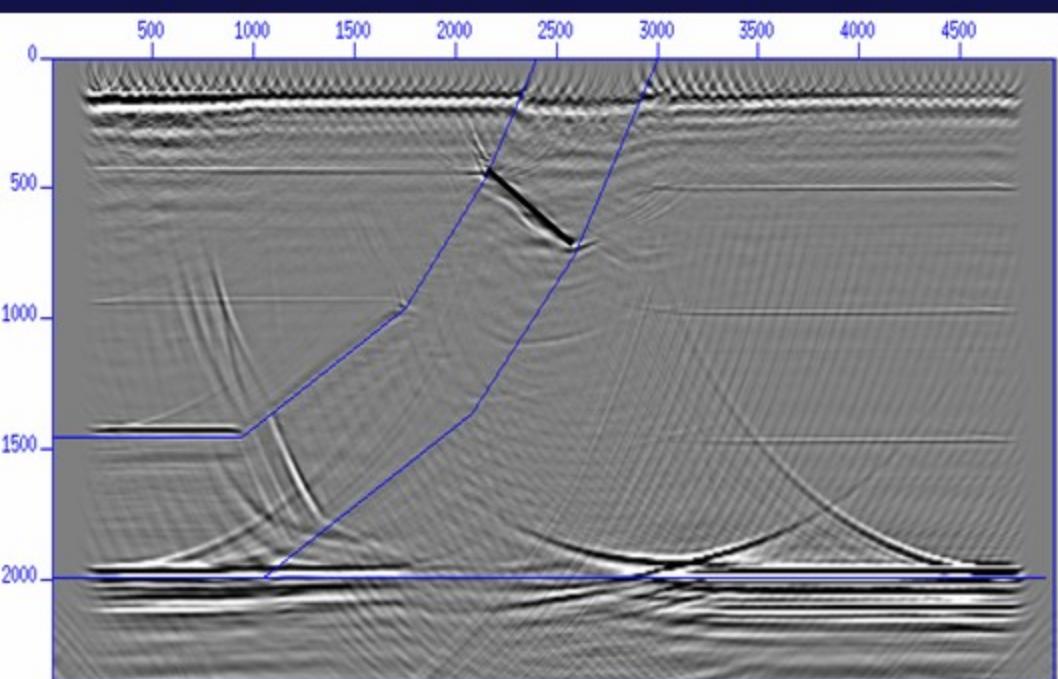


Thrust model.

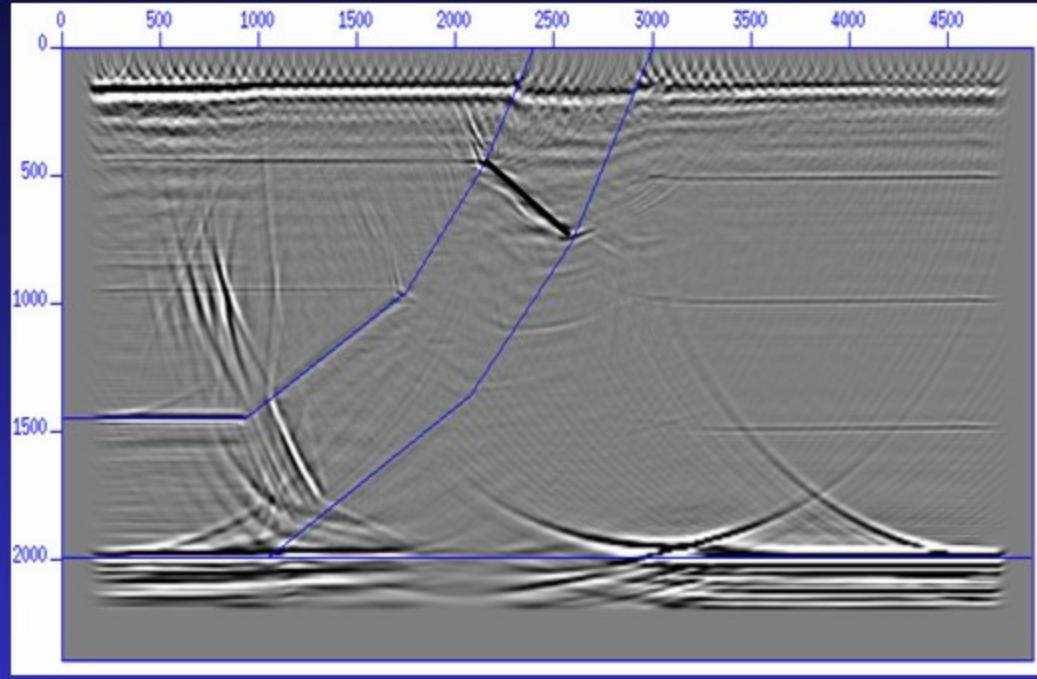
Main area of interest is the apparent pull up

Zero-offset seismic

Physical example: Poststack migration



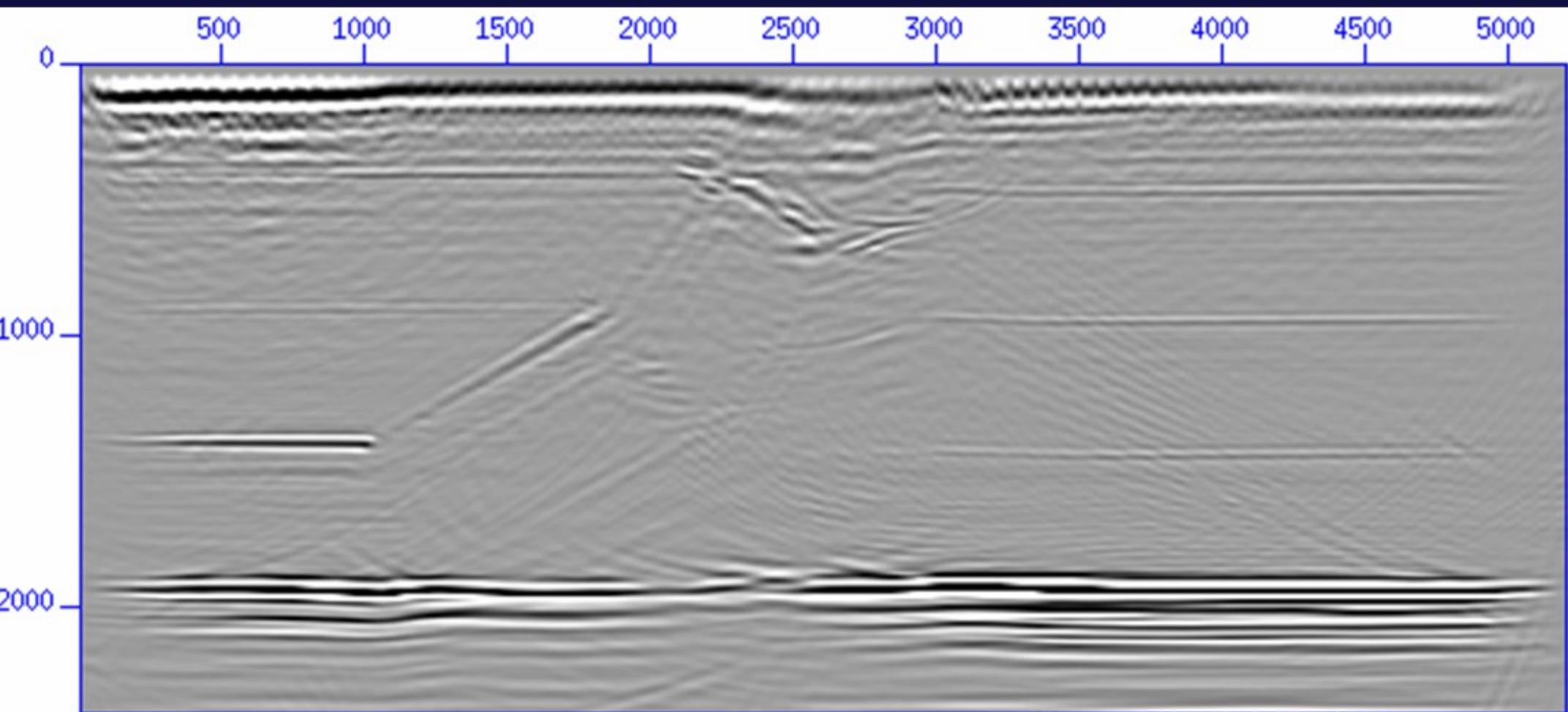
Anisotropic **PSPI** migration



Anisotropic reverse-time migration

Prestack

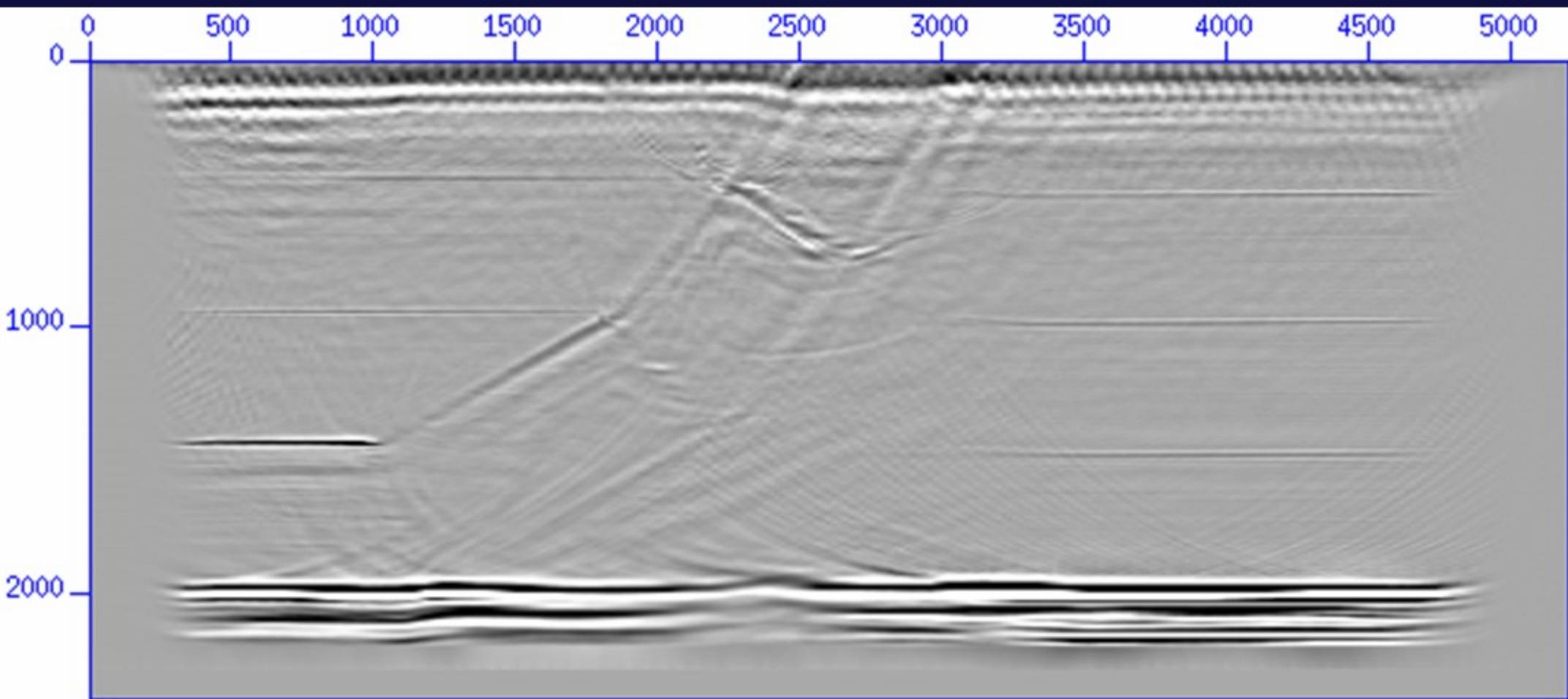
A-PSPI



Prestack **A-PSPI** migration result

Prestack

A-RT



Prestack **A-RT** migration result

Some theory: Xiang Du

P and SV phase velocity for TTI media

The VTI phase-velocity equation (Tsvankin, 1996),

$$\frac{V^2(\theta)}{V_{p0}^2} = 1 + \varepsilon \sin^2 \theta - \frac{f}{2} \pm \frac{f}{2} \sqrt{\left(1 + \frac{2\varepsilon \sin^2 \theta}{f}\right)^2 - \frac{2(\varepsilon - \delta) \sin^2 2\theta}{f}}$$

When we rotate the symmetry axis from vertical to a tilt angle ϕ , the phase velocity in the direction measured from the vertical direction becomes:

$$\frac{V^2(\theta, \phi)}{V_{p0}^2} = 1 + \varepsilon \sin^2(\theta - \phi) - \frac{f}{2} \pm \frac{f}{2} \sqrt{\left(1 + \frac{2\varepsilon \sin^2(\theta - \phi)}{f}\right)^2 - \frac{2(\varepsilon - \delta) \sin^2 2(\theta - \phi)}{f}}$$

Simplified P and SV phase velocity for TTI media

Using a weak anisotropy assumption, the familiar linear approximation can be formed:

$$\frac{V_p^2(\theta, \phi)}{V_{p0}^2} = 1 + 2\delta \sin^2(\theta - \phi) \cos^2(\theta - \phi) + 2\varepsilon \sin^4(\theta - \phi)$$

$$\frac{V_s^2(\theta, \phi)}{V_{p0}^2} = 1 - f + 2(\varepsilon - \delta) \sin^2(\theta - \phi) \cos^2(\theta - \phi)$$

Alkhalifah (2000) acoustic wave equation (By setting $f=1$) for titled angle ϕ as

$$\frac{V_p^2(\theta, \phi)}{V_{p0}^2} = \frac{1}{2} + \varepsilon \sin^2(\theta - \phi) + \frac{1}{2} \sqrt{(1 + 2\varepsilon \sin^2(\theta - \phi))^2 - 2(\varepsilon - \delta) \sin^2 2(\theta - \phi)}$$

However this equation has problems with SV wave aliasing.
Typical solution is to include an isotropic layer near the surface.

(Zhang et al., 2004)

For improved accuracy...

Retaining the quadratic terms of exact formula, we could use:

$$\frac{V_p^2(\theta, \phi)}{V_{p0}^2} = 1 + 2\delta \sin^2(\theta - \phi) \cos^2(\theta - \phi) + 2\varepsilon \sin^4(\theta - \phi)$$
$$+ \frac{4}{f}(\varepsilon - \delta)[\varepsilon \sin^2(\theta - \phi) + \delta \cos^2(\theta - \phi)] \sin^4(\theta - \phi) \cos^2(\theta - \phi)$$

$$\frac{V_s^2(\theta, \phi)}{V_{p0}^2} = 1 - f + 2(\varepsilon - \delta) \sin^2(\theta - \phi) \cos^2(\theta - \phi)$$
$$- \frac{4}{f}(\varepsilon - \delta)[\varepsilon \sin^2(\theta - \phi) + \delta \cos^2(\theta - \phi)] \sin^4(\theta - \phi) \cos^2(\theta - \phi)$$

Simplified P and SV wave equation for TTI media (Thomsen's equation with tilt)

Using the above linear relationship and Fourier transform, we can get the simplified P wave equation in wavenumber-time domain:

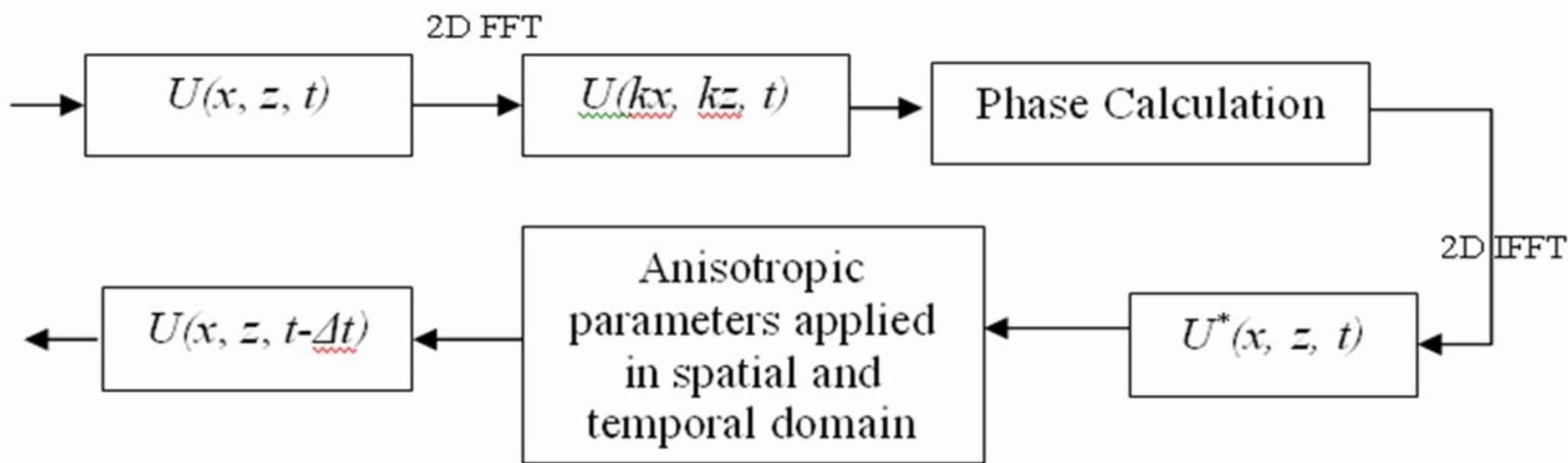
$$\frac{\partial^2 U_p(k_x, k_z, t)}{\partial t^2} = -V_{p0}^2 [k_x^2 + k_z^2 + (2\delta \sin^2 \phi \cos^2 \phi + 2\varepsilon \cos^4 \phi) \frac{k_x^4}{k_x^2 + k_z^2} + \dots + (-\delta \sin 4\phi - 4\varepsilon \sin 2\phi \sin^2 \phi) \frac{k_z^3 k_x}{k_x^2 + k_z^2})] U_p(k_x, k_z, t)$$

Simplified SV wave equation in wavenumber-time domain

$$\frac{\partial^2 U_s(k_x, k_z, t)}{\partial t^2} = -V_{s0}^2 [k_x^2 + k_z^2 + \dots + 4\varepsilon \sin 2\phi \cos^2 \phi) \frac{k_x k_z (k_x^2 - k_z^2)}{k_x^2 + k_z^2}] U_s(k_x, k_z, t)$$

$$\sigma = \left(\frac{V_{p0}}{V_{s0}}\right)^2 (\varepsilon - \delta)$$

Numerical solution: Rev. time pseudospectral method P and SV wave equations for TTI media



Computational diagram of anisotropic reverse-time scheme corresponding to P- and SV-wave equation

Anisotropic PSPI for TTI media

Considering the frequency-dispersion and phase velocity

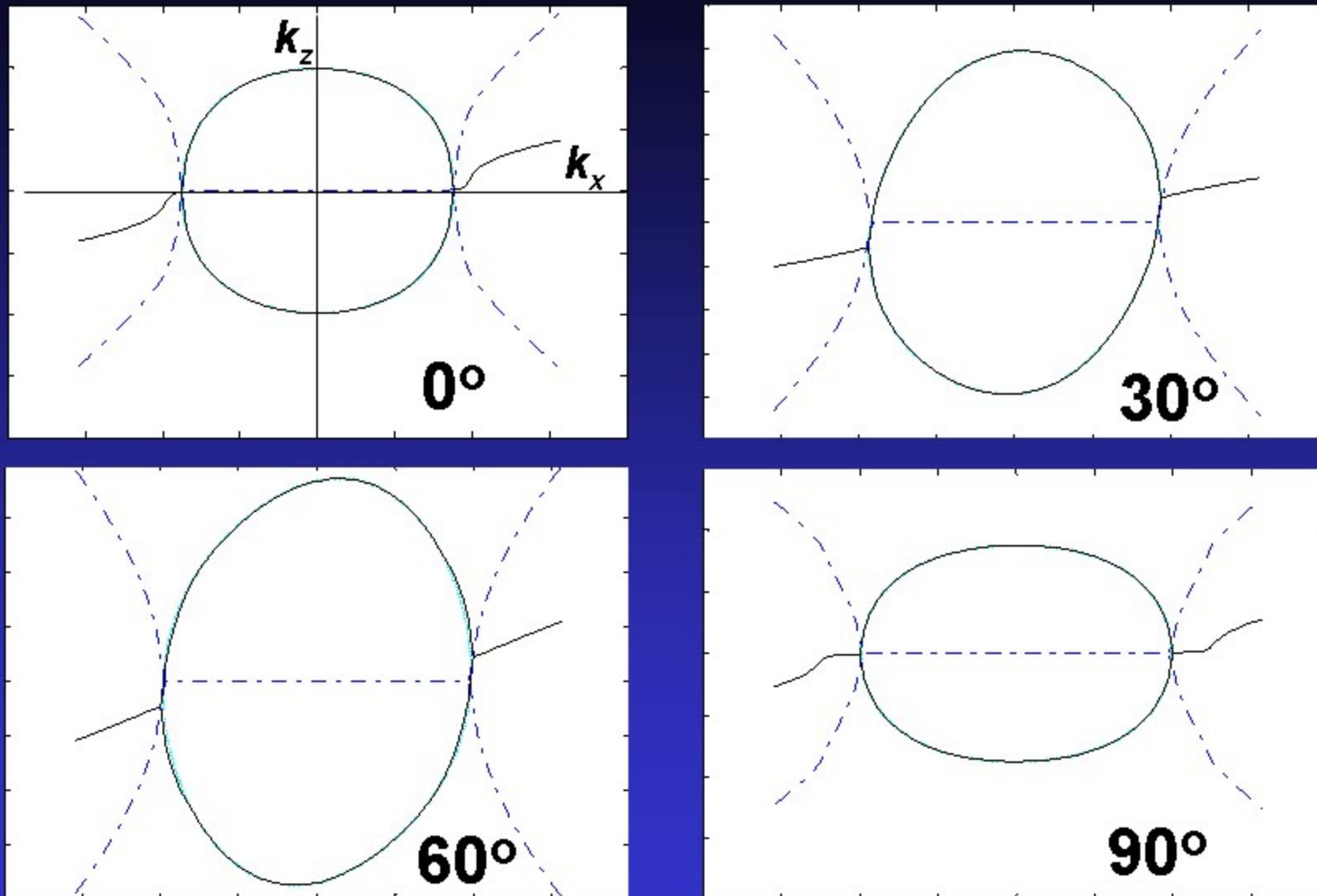
$$k_z = \pm \sqrt{\frac{\omega^2}{v^2(\theta, \phi)} - k_x^2}$$

We can write it as the quartic equation that is different from the table-driven method (1997) and polynomial approximation (2002) .

$$k_z^4 + a_3 k_z^3 + a_2 k_z^2 + a_1 k_z + a_0 = 0$$

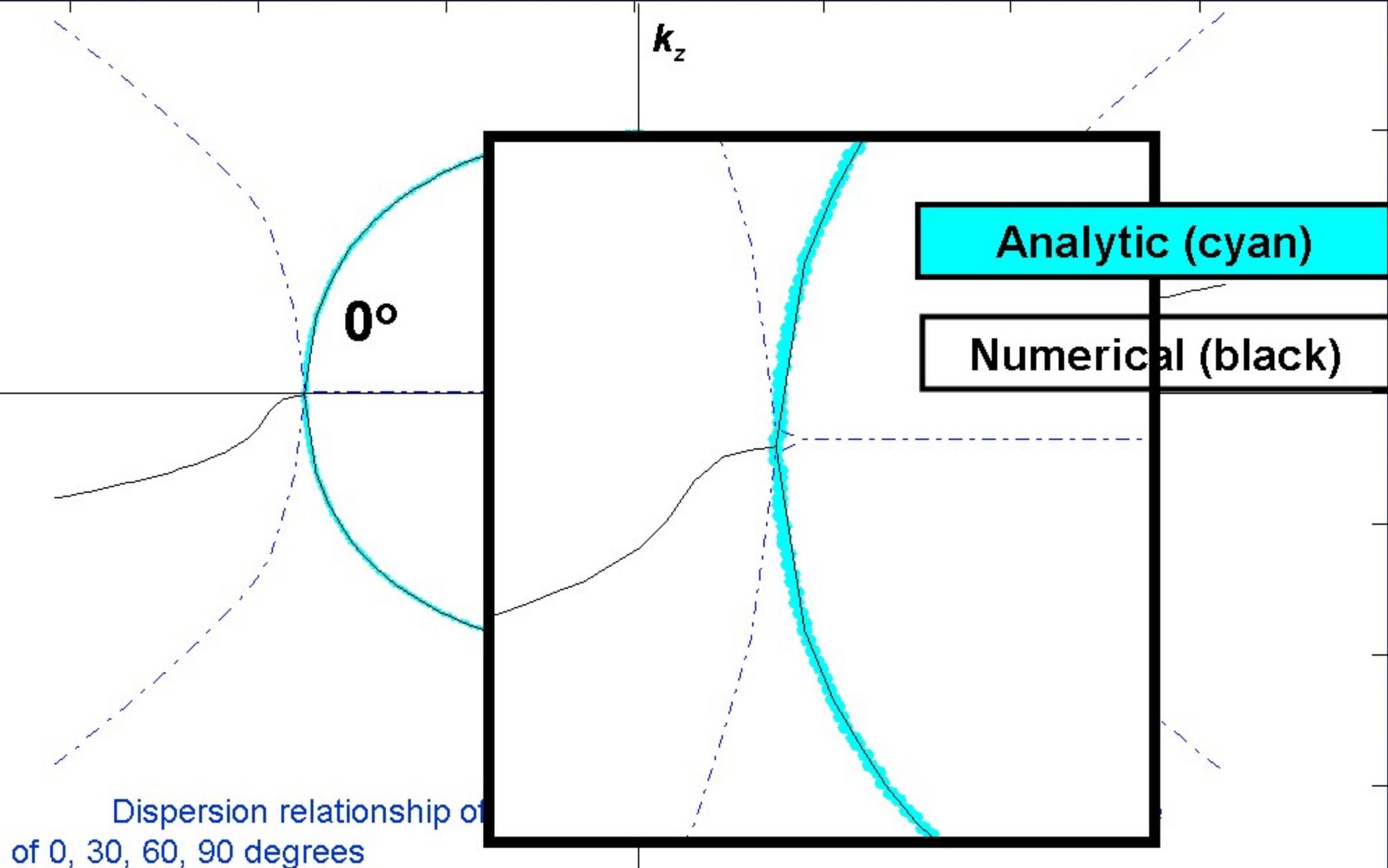
With the quartic equation, we can get the analytic solution for Kz

Dispersion relationship K_x, K_z

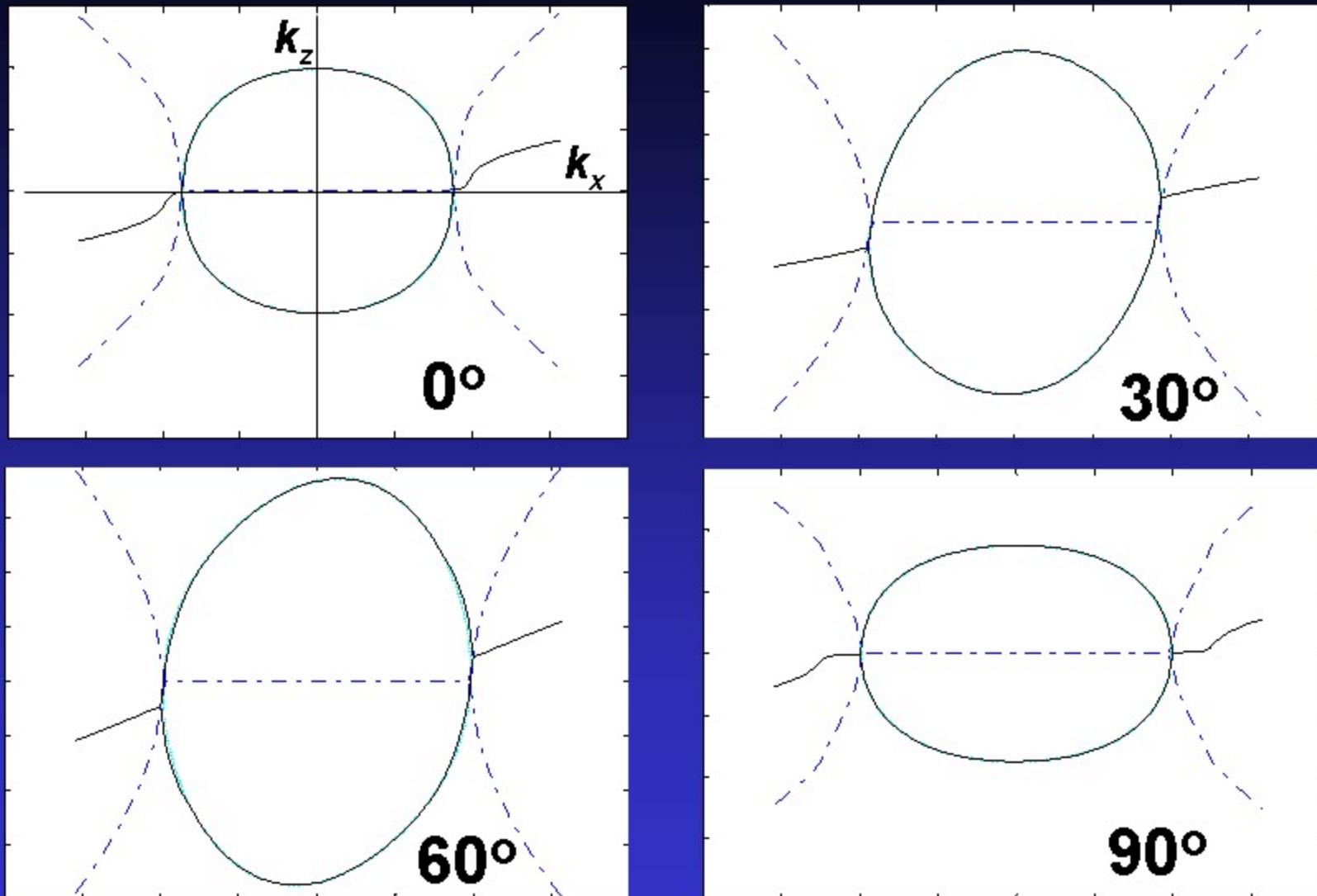


Dispersion relationship of P-wave in a TI medium with tilt angles
of 0, 30, 60, 90 degrees

Analytical and Numerical solution for k_z

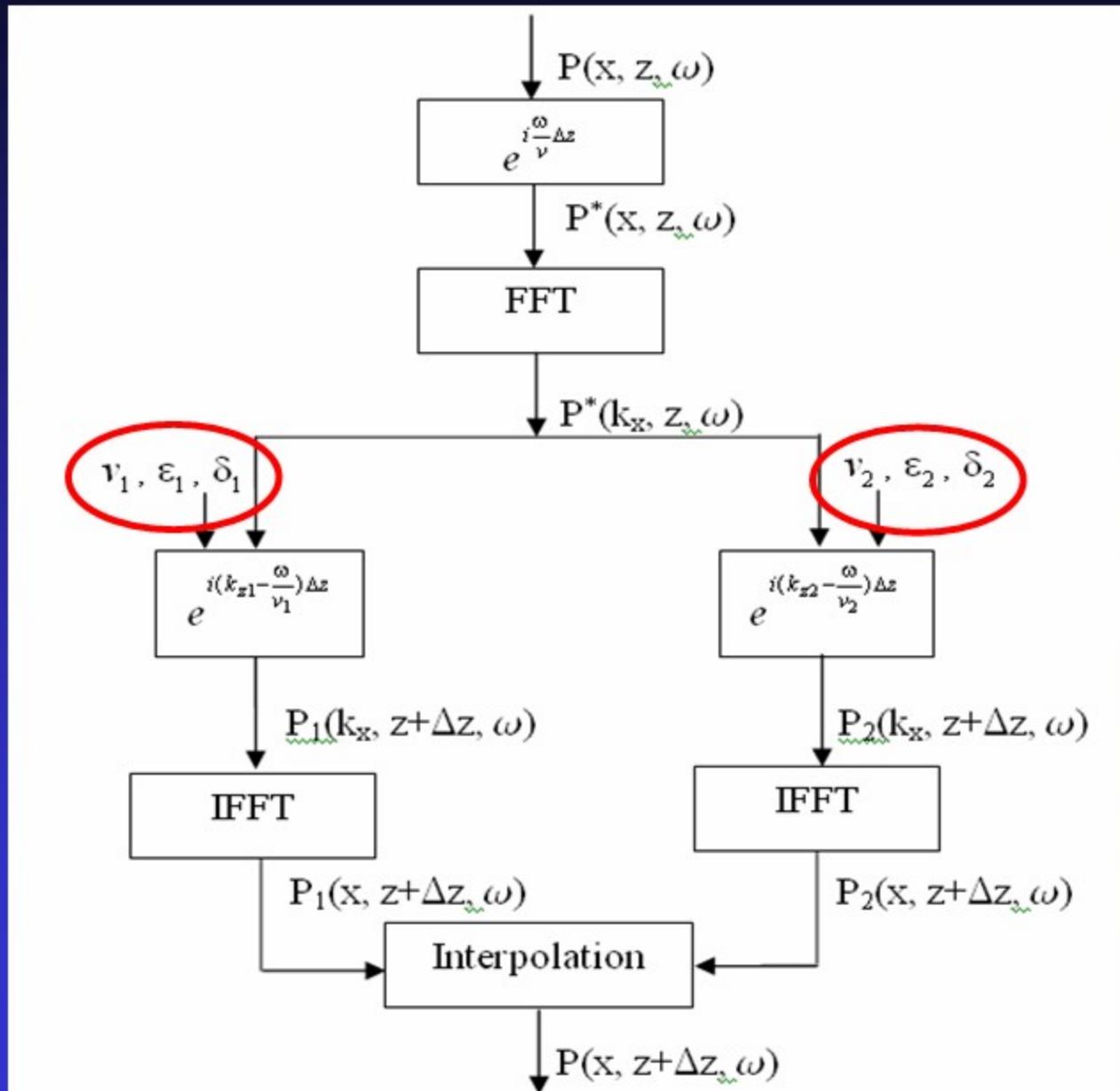


Dispersion relationship K_x, K_z



Dispersion relationship of P-wave in a TI medium with tilt angles
of 0, 30, 60, 90 degrees

Computational diagram of Anisotropic PSPI extrapolation scheme



Conclusions (Du)

- Wave-equations for P- and SV-waves were obtained for use in TTI anisotropy reverse-time migration.
- The pseudo-spectral method was easier to solve for P- and SV-wave equations implementing reverse-time migration.
- Analytical solution for K_z is employed for A-PSPI
- Anisotropic PSPI makes a good balance between efficiency and accuracy.
- Anisotropic reverse time migration shows excellent steep angle imaging.

The $\beta v \tau\tau$ stops here

